

## INFINITELY MANY SOLUTIONS FOR QUASILINEAR SCHRÖDINGER EQUATIONS UNDER BROKEN SYMMETRY SITUATION

LIANG ZHANG — XIANHUA TANG — YI CHEN

---

ABSTRACT. In this paper, we study the existence of infinitely many solutions for the quasilinear Schrödinger equations

$$\begin{cases} -\Delta u - \Delta(|u|^\alpha)|u|^{\alpha-2}u = g(x, u) + h(x, u) & \text{for } x \in \Omega, \\ u = 0 & \text{for } x \in \partial\Omega, \end{cases}$$

where  $\alpha \geq 2$ ,  $g, h \in C(\Omega \times \mathbb{R}, \mathbb{R})$ . When  $g$  is of superlinear growth at infinity in  $u$  and  $h$  is not odd in  $u$ , the existence of infinitely many solutions is proved in spite of the lack of the symmetry of this problem, by using the dual approach and Rabinowitz perturbation method. Our results generalize some known results and are new even in the symmetric situation.

### 1. Introduction and main results

Consider the following quasilinear Schrödinger equation:

$$(1.1) \quad \begin{cases} -\Delta u - \Delta(|u|^\alpha)|u|^{\alpha-2}u = g(x, u) + h(x, u) & \text{for } x \in \Omega, \\ u = 0 & \text{for } x \in \partial\Omega, \end{cases}$$

where  $\alpha \geq 2$ ,  $g, h \in C(\Omega \times \mathbb{R}, \mathbb{R})$ , and  $\Omega \subset \mathbb{R}^N$  is a bounded smooth domain.

---

2010 *Mathematics Subject Classification*. Primary: 35B05, 35B45; Secondary: 35J50.

*Key words and phrases*. Broken symmetry; dual approach; quasilinear Schrödinger equation; Rabinowitz perturbation method.

This work is partially supported by the NNSF (No.11171351, 11426211, 11571370) of China, Natural Science Foundation of Shandong Province of China (No.ZR2014AP011) and the Natural Science Foundation of Jiangsu Province of China (No.BK20140176).

In recent years, the quasilinear Schrödinger equation has been involved in several models of mathematical physics (see [8], [9], [15]). Notice that equation (1.1) is the Euler–Lagrange equation associated with the energy functional  $J: E \rightarrow \mathbb{R}$  given by

$$(1.2) \quad J(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{1}{2\alpha} \int_{\Omega} |\nabla(|u|^\alpha)|^2 dx \\ - \int_{\Omega} G(x, u) dx - \int_{\Omega} H(x, u) dx,$$

where  $E$  denotes the Hilbert space  $H_0^1(\Omega)$  equipped with the inner product

$$(u, v) = \int_{\Omega} \nabla u \nabla v dx, \quad u, v \in E.$$

By direct computation, we have

$$(1.3) \quad \frac{1}{2\alpha} \int_{\Omega} |\nabla(|u|^\alpha)|^2 dx = \frac{\alpha}{2} \int_{\Omega} |u|^{2(\alpha-1)} |\nabla u|^2 dx, \quad u \in E.$$

In view of (1.2) and (1.3),

$$J(u) = \frac{1}{2} \int_{\Omega} |\nabla u|^2 dx + \frac{\alpha}{2} \int_{\Omega} |u|^{2(\alpha-1)} |\nabla u|^2 dx - \int_{\Omega} G(x, u) dx - \int_{\Omega} H(x, u) dx,$$

for  $u \in E$ . By (1.4), the energy functional  $J$  could be naturally defined on

$$X = \left\{ u \in H_0^1(\Omega) \mid \int_{\Omega} |u|^{2(\alpha-1)} |\nabla u|^2 dx < \infty \right\},$$

which is not a vector space. So there is no suitable space on which the energy functional  $J$  is well-defined. In recent years several methods have been developed to overcome this difficulty, such as the constrained minimization (see [10]), Nehari method (see [7], [11], [18]), change of variables (dual approach) (see [1], [7], [23], [25], [26]), perturbation method (see [12], [13], [24]). Recently, Liu and Zhao [14] considered the existence of infinitely many solutions for a more general quasilinear equation

$$\begin{cases} D_j \left( \sum_{i,j=1}^N a_{ij}(x, u) D_i u \right) - \frac{1}{2} \sum_{i,j=1}^N D_s a_{ij}(x, u) D_i u D_j u + |u|^{p-2} u + f = 0 & \text{for } x \in \Omega, \\ u = 0 & \text{for } x \in \partial\Omega, \end{cases}$$

where  $D_i := \frac{\partial}{\partial x_i}$ ,  $i = 1, \dots, N$ ,  $D_s a_{ij}(x, s) = \frac{\partial}{\partial s} a_{ij}(x, s)$ . They treated the case  $f \neq 0$  as a perturbation from a symmetric equation. Under some suitable conditions, they showed the existence of infinitely many solutions for this quasilinear equation. Similar questions under symmetry breaking situation have been studied also for the problems of elliptic type, Hamiltonian systems and ordinary differential equations (see [3]–[6], [16], [19]–[22]).

But it should be noted that  $|u|^{p-2}u$  is a special form of function, which satisfies the classical condition (AR) due to Ambrosetti and Rabinowitz. The condition (AR) is a convenient hypothesis since it achieves the mountain pass geometry as well as fulfils the Palais–Smale condition, but this condition is far too restrictive. There are many functions not satisfying (AR). For example, let

$$(1.4) \quad g(x, t) = 2\alpha\theta(x)|t|^{2(\alpha-1)}t \left[ \ln(1 + |t|^{2\alpha}) + \frac{|t|^{2\alpha}}{1 + |t|^{2\alpha}} \right], \quad (x, t) \in \Omega \times \mathbb{R},$$

where  $\theta: \bar{\Omega} \rightarrow \mathbb{R}$  is a bounded continuous function with  $\inf_{x \in \bar{\Omega}} \theta(x) > 0$ . To the best of our knowledge, the question whether infinitely many solutions persist for system (1.1) with functions not satisfying (AR) with broken symmetry is unsettled. In this paper, we give a positive answer to this question. But this question is different from the case discussed in [14], those methods cannot be applied directly to obtain our results. Our main tools are based on the dual approach and Rabinowitz perturbation method introduced in [16].

**THEOREM 1.1.** *Assume that  $g$  and  $h$  satisfy the following conditions:*

(g<sub>1</sub>)  $g \in C(\Omega \times \mathbb{R}, \mathbb{R})$  and there exists  $2\alpha < p < 2^*\alpha$  if  $N \geq 3$  or  $2\alpha < p < \infty$  if  $N = 1, 2$  such that

$$|g(x, t)| \leq C_0(1 + |t|^{p-1}), \quad (x, t) \in \Omega \times \mathbb{R};$$

(g<sub>2</sub>) there exists a positive constant  $r_0 > 0$  such that

$$G(x, t) \geq 0, \quad (x, t) \in \Omega \times \mathbb{R} \quad \text{and} \quad |t| \geq r_0,$$

and

$$\lim_{|t| \rightarrow \infty} \frac{G(x, t)}{|t|^{2\alpha}} = \infty, \quad \text{a.e. } x \in \Omega;$$

where  $G(x, t) := \int_0^t g(x, s) ds$ ;

(g<sub>3</sub>) there exist constants  $C_1 > 0$  and  $\kappa > \max\{1, N/2\}$  such that

$$|G(x, t)|^\kappa \leq C_1 |t|^{2\alpha\kappa} \bar{G}(x, t), \quad (x, t) \in \Omega \times \mathbb{R}, \quad |t| \geq r_0,$$

where  $\bar{G}(x, t) := (2\alpha)^{-1}tg(x, t) - G(x, t)$ ;

(g<sub>4</sub>) there exists a positive constant  $C_2 > 0$  such that

$$\bar{G}(x, t) \geq C_2(|t|^{2\alpha} - 1), \quad (x, t) \in \Omega \times \mathbb{R};$$

(g<sub>5</sub>)  $g(x, t) = -g(x, -t)$  for  $(x, t) \in \Omega \times \mathbb{R}$ ;

(h<sub>1</sub>)  $h \in C(\Omega \times \mathbb{R}, \mathbb{R})$  and there exist constants  $C_3 > 0$  and  $1 < \sigma < 2\alpha$  such that

$$|h(x, t)| \leq C_3(1 + |t|^{\sigma-1}), \quad (x, t) \in \Omega \times \mathbb{R};$$

(h<sub>2</sub>) the constants  $p, \alpha$  and  $\sigma$  satisfy

$$\frac{2\alpha N - p(N - 2)}{N(p - 2\alpha)} > \frac{2\alpha}{2\alpha - \sigma}.$$

Then system (1.1) has an unbounded sequence of solutions.

**COROLLARY 1.2.** *Assume that  $g$  and  $h$  satisfy  $(g_1)$ – $(g_5)$ ,  $(h_1)$  and the following assumption:*

$$(h_3) \quad h(x, t) = -h(x, -t) \text{ for } (x, t) \in \Omega \times \mathbb{R}.$$

Then there exists an unbounded sequence of solutions for system (1.1).

The plan of this paper is as follows. In Section 2 we provide some preliminary materials. We prove our main results by the use of the dual approach and Rabinowitz perturbation method in Section 3. In the last section an example is given to illustrate our results.

**Notation.** Throughout the paper, we denote by  $C_n$  various positive constants, which may vary from line to line and are not essential to the proof.

## 2. Preliminaries

For any  $s \in [1, 2^*]$ ,  $L^s(\Omega)$  is the usual Lebesgue space with the norm

$$\|u\|_s := \left( \int_{\Omega} |u|^s dx \right)^{1/s},$$

and  $H_0^1(\Omega)$  is the usual Sobolev space with the norm

$$\|u\| := \left( \int_{\Omega} |\nabla u|^2 dx \right)^{1/2},$$

we denote the Hilbert space  $H_0^1(\Omega)$  by  $E$ . It is well-known that  $E$  is continuously embedded into  $L^s(\Omega)$  for  $s \in [1, 2^*]$ , i.e. there exist constants  $\tau_s > 0$  such that

$$\|u\|_s \leq \tau_s \|u\|, \quad u \in E, \quad s \in [1, 2^*].$$

Moreover,  $E \hookrightarrow L^s(\Omega)$  is compact for  $s \in [1, 2^*)$ .

It is obvious that the second order differential operator with the Dirichlet boundary condition is a selfadjoint operator, and there exists a sequence of eigenvalues (counted with multiplicity)  $\lambda_1 < \lambda_2 < \dots \rightarrow \infty$ , and the corresponding system of normalized eigenfunctions  $\{e_n : n \in \mathbb{N}\}$  forming an orthogonal basis in  $E$ . Hereafter, let  $E_n := \text{span}\{e_1, \dots, e_n\}$  and  $E_n^\perp$  be the orthogonal complement of  $E_n$  in  $E$ .

Inspired by the transformation initially introduced in [9], the function  $f$  can be defined by

$$\begin{aligned} f'(t) &= \frac{1}{\sqrt{1 + \alpha|f(t)|^{2(\alpha-1)}}} && \text{for } t \in [0, +\infty), \\ f(-t) &= -f(t) && \text{for } t \in (-\infty, 0]. \end{aligned}$$

Next we collect some useful properties of the function  $f: \mathbb{R} \rightarrow \mathbb{R}$ , which will be used frequently in the sequel of the paper. Proofs can be found in [1].

LEMMA 2.1. *The function  $f$  and its derivative have the following properties:*

- (f<sub>1</sub>)  $f$  is a uniquely defined  $C^\infty$  function and it is invertible;
- (f<sub>2</sub>)  $0 < f'(t) \leq 1$  and  $|f(t)| \leq |t|$ , for all  $t \in \mathbb{R}$ ;
- (f<sub>3</sub>)  $\lim_{t \rightarrow 0} |f(t)|/|t| = 1$  and  $\lim_{t \rightarrow \infty} |f(t)|^\alpha/|t| = \sqrt{\alpha}$ ;
- (f<sub>4</sub>) there exists a positive constant  $C_0$  such that

$$|f(t)|^{\alpha-1} f'(t) \leq C_0, \quad \text{for all } t \in \mathbb{R};$$

- (f<sub>5</sub>)  $f''(t)f(t) = (\alpha - 1)(f'(t))^2((f'(t))^2 - 1)$ , for all  $t \in \mathbb{R}$ .

Therefore, after change of variables, we obtain the following functional:

$$(2.1) \quad I(v) := J(f(v)) = \frac{1}{2} \int_{\Omega} |\nabla v|^2 dx - \int_{\Omega} G(x, f(v)) dx - \int_{\Omega} H(x, f(v)) dx.$$

Moreover, for any  $v, w \in E$ ,

$$(2.2) \quad \langle I'(v), w \rangle = (v, w) - \int_{\Omega} g(x, f(v)) f'(v) w dx - \int_{\Omega} h(x, f(v)) f'(v) w dx.$$

By a standard argument which is similar to Lemma 2.6 and Remark 2.7 in [1], if  $v \in E$  is a critical point of the functional  $I$ , then  $u = f(v) \in E$  and  $u$  is a weak solution of (1.1).

In order to define a suitable modified functional, we prove the following lemma.

LEMMA 2.2. *Under the hypotheses of Theorem 1.1, there exists a positive constant  $A$  depending on  $\alpha$  such that if  $v$  is a critical point of  $I$ ,*

$$(2.3) \quad \int_{\Omega} |f(v)|^{2\alpha} dx \leq A(I^2(v) + 1)^{1/2}.$$

PROOF. Since  $v$  is a critical point of  $I$ , by (f<sub>5</sub>), (g<sub>4</sub>), (h<sub>1</sub>), (2.1) and (2.2),

$$(2.4) \quad I(v) - \frac{1}{2\alpha} \left\langle I'(v), \frac{f(v)}{f'(v)} \right\rangle > \int_{\Omega} \bar{G}(x, f(v)) dx - C_4 \left( \int_{\Omega} |f(v)|^\sigma dx + 1 \right) > C_5 \int_{\Omega} |f(v)|^{2\alpha} dx - C_6.$$

Then (2.3) follows from (2.4) and the Young inequality. □

Next we introduce a cut-off function  $\zeta \in C^\infty(\mathbb{R}, \mathbb{R})$  such that

$$(2.5) \quad \begin{cases} \zeta(t) = 1 & \text{for } t \in (-\infty, 1], \\ 0 \leq \zeta(t) \leq 1 & \text{for } t \in (1, 2), \\ \zeta(t) = 0 & \text{for } t \in [2, \infty), \\ |\zeta'(t)| \leq 2 & \text{for } t \in \mathbb{R}. \end{cases}$$

With the help of this cut-off function, define

$$(2.6) \quad P(v) = 2A(I^2(v) + 1)^{1/2}, \quad \phi(v) = \zeta \left( P^{-1}(v) \int_{\Omega} |f(v)|^{2\alpha} dx \right).$$

If  $v$  is a critical point of  $I$ , by (2.3), (2.5) and (2.6),  $\phi(v) = 1$ . Set

$$(2.7) \quad \bar{I}(v) = \frac{1}{2} \|v\|^2 - \int_{\Omega} G(x, f(v)) dx - \phi(v) \int_{\Omega} H(x, f(v)) dx.$$

Since  $\zeta$  is a smooth function, we have  $\bar{I} \in C^1(E, \mathbb{R})$  and

$$(2.8) \quad \begin{aligned} \langle \bar{I}'(v), w \rangle &= (v, w) - \int_{\Omega} g(x, f(v)) f'(v) w dx \\ &\quad - \phi(v) \int_{\Omega} h(x, f(v)) f'(v) w dx - \langle \phi'(v), w \rangle \int_{\Omega} H(x, f(v)) dx, \end{aligned}$$

for  $v, w \in E$ . It is obvious that  $I(v) = \bar{I}(v)$  if  $v$  is a critical point of  $I$ .

LEMMA 2.3. *Assume all the hypotheses of Theorem 1.1 hold. Then:*

(H<sub>1</sub>) *there is a positive constant  $C_7$  such that*

$$|\bar{I}(v) - \bar{I}(-v)| \leq C_7(|\bar{I}(v)|^{\sigma/2\alpha} + 1), \quad v \in E;$$

(H<sub>2</sub>) *there exists a positive constant  $M_1$  such that if  $\bar{I}(v) \geq M_1$  and  $\bar{I}'(v) = 0$ , then  $\bar{I}(v) = I(v)$  and  $I'(v) = 0$ ;*

(H<sub>3</sub>) *there exists a positive constant  $M_2 \geq M_1$  such that for any  $c > M_2$ , then  $\bar{I}$  satisfies the  $(C)_c$  condition at  $c$ .*

PROOF. If  $v \in \text{supp } \phi$ , by (2.5) and (2.6),

$$(2.9) \quad \int_{\Omega} |f(v)|^{2\alpha} dx \leq 4A(I^2(v) + 1)^{1/2} \leq 4A(|I(v)| + 1).$$

By (h<sub>1</sub>) and direct computation, we have

$$(2.10) \quad \left| \int_{\Omega} H(x, f(v)) dx \right| \leq C_8 \left( \int_{\Omega} |f(v)|^{2\alpha} dx + 1 \right)^{\sigma/2\alpha}.$$

It follows from (2.9) and (2.10) that

$$(2.11) \quad \left| \int_{\Omega} H(x, f(v)) dx \right| \leq C_9(|I(v)|^{\sigma/2\alpha} + 1).$$

In view of (2.1), (2.7) and (2.11),

$$(2.12) \quad |I(v)| \leq |\bar{I}(v)| + \left| \int_{\Omega} H(x, f(v)) dx \right| \leq |\bar{I}(v)| + 2C_9(|I(v)|^{\sigma/2\alpha} + 1).$$

In combination with (h<sub>1</sub>) and (2.12),

$$(2.13) \quad |I(v)| \leq C_{10}(|\bar{I}(v)| + 1).$$

It follows from (2.11) and (2.13) that

$$(2.14) \quad \left| \int_{\Omega} H(x, f(v)) dx \right| \leq C_{11}(|\bar{I}(v)|^{\sigma/2\alpha} + 1), \quad v \in \text{supp } \phi.$$

By a similar estimate, we also have

$$(2.15) \quad \left| \int_{\Omega} H(x, f(-v)) dx \right| \leq C_{12}(|\bar{I}(v)|^{\sigma/2\alpha} + 1), \quad -v \in \text{supp } \phi.$$

It follows from (g<sub>5</sub>), (2.7), (2.14) and (2.15) that (H<sub>1</sub>) holds.

To prove (H<sub>2</sub>), it suffices to show that  $v$  is a critical point of  $\bar{I}$  with  $\bar{I}(v) \geq M_1$ , then

$$(2.16) \quad P^{-1}(v) \int_{\Omega} |f(v)|^{2\alpha} dx < 1.$$

Next we show that (2.16) holds. It follows from (2.8) that

$$(2.17) \quad \begin{aligned} \left\langle \bar{I}'(v), \frac{f(v)}{f'(v)} \right\rangle &= \alpha \|v\|^2 - (\alpha - 1) \int_{\Omega} (f'(v))^2 |\nabla v|^2 dx \\ &\quad - \int_{\Omega} g(x, f(v)) f(v) dx - \phi(v) \int_{\Omega} h(x, f(v)) f(v) dx \\ &\quad - \left\langle \phi'(v), \frac{f(v)}{f'(v)} \right\rangle \int_{\Omega} H(x, f(v)) dx, \end{aligned}$$

where

$$\begin{aligned} \langle \phi'(v), \omega \rangle &= \zeta'(\theta(v)) P^{-2}(v) \\ &\quad \cdot \left[ 2\alpha P(v) \int_{\Omega} |f(v)|^{2(\alpha-1)} f(v) f'(v) \omega dx - (2A)^2 \theta(v) I(v) \langle I'(v), \omega \rangle \right], \end{aligned}$$

and

$$(2.18) \quad \theta(v) := P^{-1}(v) \int_{\Omega} |f(v)|^{2\alpha} dx.$$

If  $v \notin \text{supp } \phi$ ,  $\phi(v) = \phi'(v) = 0$ , then (2.17) reduces to

$$(2.19) \quad \begin{aligned} \left\langle \bar{I}'(v), \frac{f(v)}{f'(v)} \right\rangle &= \alpha \|v\|^2 - (\alpha - 1) \int_{\Omega} (f'(v))^2 |\nabla v|^2 dx \\ &\quad - \int_{\Omega} g(x, f(v)) f(v) dx. \end{aligned}$$

Moreover, if  $v$  is a critical point of  $\bar{I}$ , by (g<sub>3</sub>), (2.1), (2.8) and (2.19),

$$(2.20) \quad \begin{aligned} I(v) - \frac{1}{2\alpha} \left\langle \bar{I}'(v), \frac{f(v)}{f'(v)} \right\rangle &> \frac{\alpha - 1}{2\alpha} \int_{\Omega} (f'(v))^2 |\nabla v|^2 dx \\ &\quad + \int_{\Omega} \bar{G}(x, f(v)) dx - \int_{\Omega} H(x, f(v)) dx. \end{aligned}$$

In view of (2.4) and (2.20), (2.16) holds. If  $v \in \text{supp } \phi$ , we regroup terms in (2.17) yielding

$$(2.21) \quad \begin{aligned} \left\langle \bar{I}'(v), \frac{f(v)}{f'(v)} \right\rangle &= (1 + K_1(v)) \left[ \alpha \|v\|^2 - (\alpha - 1) \int_{\Omega} (f'(v))^2 |\nabla v|^2 dx \right] \\ &\quad - K_2(v) \int_{\Omega} H(x, f(v)) dx - (1 + K_1(v)) \int_{\Omega} g(x, f(v)) f(v) dx \\ &\quad - (\phi(v) + K_1(v)) \int_{\Omega} h(x, f(v)) f(v) dx, \end{aligned}$$

where

$$(2.22) \quad K_1(v) := (2A)^2 \zeta'(\theta(v)) P^{-2}(v) \theta(v) I(v) \int_{\Omega} H(x, f(v)) dx,$$

$$(2.23) \quad K_2(v) := 2\alpha \zeta'(\theta(v)) \theta(v).$$

Next we prove that

$$(2.24) \quad K_1(v) \rightarrow 0, \quad \text{as } M_1 \rightarrow \infty.$$

By (2.5), (2.6), (2.11) and (2.22),

$$(2.25) \quad |K_1(v)| \leq 8C_9(|I(v)|^{\sigma/2\alpha} + 1)|I(v)|^{-1}.$$

In combination with (2.1) and (2.7),

$$(2.26) \quad I(v) \geq \bar{I}(v) - \left| \int_{\Omega} H(x, f(v)) dx \right|.$$

In view of (2.11) and (2.26),

$$(2.27) \quad I(v) + C_9|I(v)|^{\sigma/2\alpha} \geq \bar{I}(v) - C_9 \geq \frac{M_1}{2},$$

for  $M_1$  large enough. If  $I(v) \leq 0$ , by (2.27) and the Young inequality,

$$(2.28) \quad \frac{(2\alpha - \sigma)C_9^{2\alpha/(2\alpha - \sigma)}}{2\alpha} + \frac{\sigma|I(v)|}{2\alpha} \geq \frac{M_1}{2} + |I(v)|.$$

But the above inequality is impossible if  $M_1$  is large enough, e.g.  $M_1 \geq (2\alpha - \sigma) \cdot C_9^{2\alpha/(2\alpha - \sigma)}/\alpha$ . Therefore  $I(v) > 0$ . Hence it follows from (2.27) that

$$I(v) > \frac{M_1}{4} \quad \text{or} \quad I(v) > \left( \frac{M_1}{2C_9} \right)^{2\alpha/\sigma},$$

which implies that

$$(2.29) \quad I(v) \rightarrow +\infty, \quad \text{as } M_1 \rightarrow \infty,$$

which together with (2.25) shows that (2.24) holds. Moreover, it follows from (2.5), (2.6) and (2.23) that  $|K_2(v)| \leq 8\alpha$ ,  $v \in E$ .

If  $v$  is a critical point of  $\bar{I}$  and  $M_1$  is large enough such that  $|K_1(v)| < 1/2$ , it follows from (2.1) and (2.21) that

$$(2.30) \quad I(v) - \frac{1}{2\alpha(1 + K_1(v))} \left\langle \bar{I}'(v), \frac{f(v)}{f'(v)} \right\rangle > \frac{\alpha - 1}{2\alpha} \int_{\Omega} (f'(v))^2 |\nabla v|^2 dx \\ + \int_{\Omega} \bar{G}(x, f(v)) dx - 8\alpha \int_{\Omega} [|h(x, f(v))f(v)| + |H(x, f(v))|] dx.$$

In view of (h<sub>1</sub>), (2.4) and (2.30), we can replace  $A$  by a larger constant but smaller than  $2A$  in (2.3), then (2.16) holds.

To prove (H<sub>3</sub>), first we show that there exists  $M_2 > M_1$  such that if  $\{v_n\}_{n \in \mathbb{N}} \subset E$  is a sequence such that

$$(2.31) \quad \bar{I}(v_n) \rightarrow c \quad \text{and} \quad \|\bar{I}'(v_n)\|(1 + \|v_n\|) \rightarrow 0,$$

then  $(v_n)$  is bounded. To prove the boundedness of  $\{v_n\}$ , arguing by contradiction, suppose that  $\|v_n\| \rightarrow \infty$ . Let  $w_n = v_n/\|v_n\|$ . Then  $\|w_n\| = 1$  and  $\|v_n\|_s \leq \tau_s \|v_n\| = \tau_s$  for  $2 \leq s < 2^*$ . Passing to a subsequence, we assume that  $w_n \rightharpoonup w$  in  $E$ , then  $w_n \rightarrow w$  in  $L^s(\Omega)$ ,  $2 \leq s < 2^*$ . For  $0 \leq a < b$ , let

$$(2.32) \quad \Omega_n(a, b) = \{x \in \Omega : a \leq |f(v_n(x))| < b\}.$$

In view of  $(g_4)$ ,  $(h_1)$ , (2.7), (2.8) and (2.21), for  $n$  large enough

$$(2.33) \quad \begin{aligned} c + 1 &\geq \bar{I}(v_n) - \frac{1}{2\alpha(1 + K_1(v_n))} \left\langle \bar{I}'(v_n), \frac{f(v_n)}{f'(v_n)} \right\rangle \\ &> \frac{1}{2} \int_{\Omega_n(r_0, \infty)} \bar{G}(x, f(v_n)) \, dx - C_{13}, \end{aligned}$$

where  $C_{13}$  is a positive constant independent of  $n$ . By  $(f_3)$  and  $(h_1)$ ,

$$(2.34) \quad \int_{\Omega} |H(x, f(v))| \, dx \leq C_{14}(\|v\|^{\sigma/\alpha} + 1).$$

It follows from (2.7), (2.31) and (2.34) that

$$(2.35) \quad \limsup_{n \rightarrow \infty} \int_{\Omega} \frac{|G(x, f(v_n))|}{\|v_n\|^2} \, dx \geq \frac{1}{2}.$$

If  $w = 0$ , then  $w_n \rightarrow 0$  in  $L^s(\Omega)$ ,  $2 \leq s < 2^*$ ,  $w_n \rightarrow 0$  almost everywhere on  $\Omega$ . Set  $\kappa' = \kappa/(\kappa - 1)$ . Since  $\kappa > \max\{1, N/2\}$ , then  $2\kappa' \in (2, 2^*)$ . It follows from  $(g_1)$  and (2.33) that

$$(2.36) \quad \begin{aligned} &\int_{\Omega_n(r_0, \infty)} \frac{|G(x, f(v_n))|}{|v_n|^2} |w_n|^2 \, dx \\ &\leq \left[ \int_{\Omega_n(r_0, \infty)} \left( \frac{|G(x, f(v_n))|}{f^{2\alpha}(v_n)} \right)^{\kappa} \, dx \right]^{1/\kappa} \left[ \int_{\Omega_n(r_0, \infty)} |w_n|^{2\kappa'} \, dx \right]^{1/\kappa'} \\ &\leq C_{15} \left[ \int_{\Omega_n(r_0, \infty)} \bar{G}(x, f(v_n)) \, dx \right]^{1/\kappa} \left( \int_{\Omega_n(r_0, \infty)} |w_n|^{2\kappa'} \, dx \right)^{1/\kappa'} \\ &\leq C_{15} \left( \int_{\Omega} |w_n|^{2\kappa'} \, dx \right)^{1/\kappa'} \rightarrow 0. \end{aligned}$$

Combining  $(g_1)$  and (2.36), we have

$$\begin{aligned} \int_{\Omega} \frac{|G(x, f(v_n))|}{\|v_n\|^2} \, dx &= \int_{\Omega_n(0, r_0)} \frac{|G(x, f(v_n))|}{\|v_n\|^2} \, dx \\ &\quad + \int_{\Omega_n(r_0, \infty)} \frac{|G(x, f(v_n))|}{|v_n|^2} |w_n|^2 \, dx \rightarrow 0, \end{aligned}$$

which contradicts (2.35).

Set  $\Pi = \{x \in \Omega : w(x) \neq 0\}$ . If  $w \neq 0$ , then  $\text{meas}(\Pi) > 0$ . Moreover,

$$(2.37) \quad \lim_{n \rightarrow \infty} |v_n(x)| = \infty, \quad \text{a.e. } x \in \Pi.$$

It follows from (f<sub>3</sub>) and (2.32) that  $\Pi \subset \Omega_n(r_0, \infty)$  for large  $n \in \mathbb{N}$ . By (g<sub>1</sub>), (h<sub>1</sub>), (f<sub>3</sub>), (2.34), (2.37) and Fatou's lemma,

$$\begin{aligned}
0 &= \limsup_{n \rightarrow \infty} \frac{\bar{I}(v_n)}{\|v_n\|^2} = \limsup_{n \rightarrow \infty} \left[ \frac{1}{2} - \int_{\Omega} \frac{G(x, f(v_n))}{\|v_n\|^2} dx \right] \\
&= \limsup_{n \rightarrow \infty} \left[ \frac{1}{2} - \int_{\Omega_n(0, r_0)} \frac{G(x, f(v_n))}{\|v_n\|^2} dx - \int_{\Omega_n(r_0, \infty)} \frac{G(x, f(v_n))}{|v_n|^2} |w_n|^2 dx \right] \\
&\leq \limsup_{n \rightarrow \infty} \left[ \frac{1}{2} + C_0(r_0 + r_0^p) \text{meas}(\Omega) \|v_n\|^{-2} - \int_{\Omega_n(r_0, \infty)} \frac{G(x, f(v_n))}{|v_n|^2} |w_n|^2 dx \right] \\
&\leq \frac{1}{2} - \liminf_{n \rightarrow \infty} \int_{\Omega_n(r_0, \infty)} \frac{G(x, f(v_n))}{f^{2\alpha}(v_n)} \cdot \frac{f^{2\alpha}(v_n)}{|v_n|^2} |w_n|^2 dx \\
&= \frac{1}{2} - \liminf_{n \rightarrow \infty} \int_{\Omega} \frac{G(x, f(v_n))}{f^{2\alpha}(v_n)} \cdot \frac{f^{2\alpha}(v_n)}{|v_n|^2} |w_n|^2 [\chi_{\Omega_n(r_0, \infty)}(x)] dx \\
&\leq \frac{1}{2} - \int_{\Omega} \liminf_{n \rightarrow \infty} \frac{G(x, f(v_n))}{f^{2\alpha}(v_n)} \cdot \frac{f^{2\alpha}(v_n)}{|v_n|^2} |w_n|^2 [\chi_{\Omega_n(r_0, \infty)}(x)] dx = -\infty,
\end{aligned}$$

which is a contradiction. Thus  $\{v_n\}$  is bounded in  $E$ .

Since  $E$  is a reflexive space, passing to a subsequence, also denoted by  $\{v_n\}$ , it can be assumed that  $v_n \rightharpoonup v_0$ ,  $n \rightarrow \infty$ . By (f<sub>3</sub>) in Lemma 2.1, there exists a positive constant  $M_3$  such that

$$(2.38) \quad |f(t)| \leq C_{16}|t|^{1/\alpha}, \quad |t| \geq M_3.$$

For any  $v, w \in E$ , by (f<sub>2</sub>), (f<sub>4</sub>), (2.38) and the Hölder inequality,

$$\begin{aligned}
(2.39) \quad &\int_{\Omega} |f(v)|^{p-1} f'(v) |w| dx \\
&= \int_{\Omega_0} |f(v)|^{p-1} f'(v) |w| dx + \int_{\Omega \setminus \Omega_0} |f(v)|^{p-1} f'(v) |w| dx \\
&\leq C_0 C_{16} \int_{\Omega_0} |v|^{(p-\alpha)/\alpha} |w| dx + \int_{\Omega \setminus \Omega_0} |v|^{p-1} |w| dx \\
&\leq C_0 C_{16} \|v\|_{p/\alpha}^{(p-\alpha)/\alpha} \|w\|_{p/\alpha} + M_3^{p-1} \|w\|_1,
\end{aligned}$$

where  $\Omega_0 := \{x \in \Omega : |v(x)| \geq M_3\}$ . By (2.8), we have

$$\begin{aligned}
\langle \bar{I}'(v_n), v_n - v_0 \rangle &= (v_n, v_n - v_0) - \int_{\Omega} g(x, f(v_n)) f'(v_n) (v_n - v_0) dx \\
&\quad - \langle \phi'(v_n), v_n - v_0 \rangle \int_{\Omega} H(x, f(v_n)) dx \\
&\quad - \phi(v_n) \int_{\Omega} h(x, f(v_n)) f'(v_n) (v_n - v_0) dx,
\end{aligned}$$

where

$$\begin{aligned}
 (2.40) \quad & \langle \phi'(v_n), v_n - v_0 \rangle \int_{\Omega} H(x, f(v_n)) dx \\
 & = 2\alpha \zeta'(\theta(v_n)) P^{-1}(v_n) \int_{\Omega} |f(v_n)|^{2(\alpha-1)} f(v_n) f'(v_n) (v_n - v_0) dx \\
 & \quad \cdot \int_{\Omega} H(x, f(v_n)) dx - K_1(v_n) \langle I'(v_n), v_n - v_0 \rangle.
 \end{aligned}$$

If  $v_n \notin \text{supp } \phi$ ,  $\phi(v_n) = \phi'(v_n) = 0$ . Then

$$\langle \bar{I}(v_n), v_n - v_0 \rangle = (v_n, v_n - v_0) - \int_{\Omega} g(x, f(v_n)) f'(v_n) (v_n - v_0) dx.$$

Otherwise,  $v_n \in \text{supp } \phi$ , in combination with (2.6) and (2.11), we have

$$(2.41) \quad \left| P^{-1}(v_n) \int_{\Omega} H(x, f(v_n)) dx \right| \leq (2A)^{-1} C_9 (|I(v_n)|^{\sigma/2\alpha} + 1) |I(v_n)|^{-1}.$$

When  $M_2$  is large enough, in view of (2.24), (2.29) and (2.41),

$$(2.42) \quad \left| P^{-1}(v_n) \int_{\Omega} H(x, f(v_n)) dx \right| \leq \frac{1}{16}, \quad |K_1(v_n)| \leq \frac{1}{16}.$$

It follows from (g<sub>1</sub>), (h<sub>1</sub>) and (2.39) that

$$(2.43) \quad \int_{\Omega} |f(v_n)|^{2\alpha-1} f'(v_n) (v_n - v_0) dx \rightarrow 0,$$

and

$$(2.44) \quad \int_{\Omega} g(x, f(v_n)) f'(v_n) (v_n - v_0) dx \rightarrow 0,$$

$$(2.45) \quad \int_{\Omega} h(x, f(v_n)) f'(v_n) (v_n - v_0) dx \rightarrow 0.$$

In combination with (2.31), (2.40), (2.42)–(2.45),  $v_n \rightarrow v_0$ ,  $n \rightarrow \infty$ . □

LEMMA 2.4. *Under assumptions (g<sub>1</sub>), (g<sub>3</sub>) and (h<sub>1</sub>), for any finite dimensional subspace  $\tilde{E} \subset E$ ,*

$$\bar{I}(v) \rightarrow -\infty, \quad \|v\| \rightarrow \infty, \quad v \in \tilde{E}.$$

PROOF. Arguing indirectly, assume that for some sequence  $\{v_n\} \subset \tilde{E}$  with  $\|v_n\| \rightarrow \infty$ , there is  $M > 0$  such that  $\bar{I}(v_n) \geq -M$  for all  $n \in \mathbb{N}$ . Set  $w_n = v_n / \|v_n\|$ , then  $\|w_n\| = 1$ . Passing to a subsequence, we can assume that  $w_n \rightharpoonup w$  in  $E$ . Since  $\tilde{E}$  is a finite dimensional space, then  $w_n \rightarrow w \in \tilde{E}$  and  $\|w\| = 1$ . Hence, we can conclude a contradiction by a similar fashion as in Lemma 2.3. □

**3. Construction of minimax sequences and proof of Theorem 1.1**

By Lemma 2.4, there exists a strictly increasing sequence of numbers  $R_n$  such that  $\bar{I}(v) \leq 0$  for  $v \in E_n \setminus B_{R_n}$ , where  $B_{R_n}$  denotes the open ball of radius  $R_n$  centred at 0 in  $E$ , and  $\bar{B}_{R_n}$  denotes the closure of  $B_{R_n}$  in  $E$ . Next we introduce some continuous maps in  $E$ . Set

$$(3.1) \quad \Gamma_n = \{ \zeta \in C(D_n, E) : \zeta \text{ is odd and } \zeta = \text{id on } \partial B_{R_n} \cap E_n \},$$

where  $D_n := \bar{B}_{R_n} \cap E_n$ , and

$$(3.2) \quad \Lambda_n := \{ \gamma \in C(U_n, E) : \gamma|_{D_n} \in \Gamma_n \text{ and } \gamma = \text{id} \\ \text{for } v \in Q_n := (\partial B_{R_{n+1}} \cap E_{n+1}) \cup ((B_{R_{n+1}} \setminus \bar{B}_{R_n}) \cap E_n) \},$$

where

$$(3.3) \quad U_n := \{ v = te_{n+1} + \omega : t \in [0, R_{n+1}], \omega \in \bar{B}_{R_{n+1}} \cap E_n, \|v\| \leq R_{n+1} \}.$$

With the help of these continuous maps, we define two sequences of minimax values

$$(3.4) \quad b_n = \inf_{\zeta \in \Gamma_n} \max_{v \in D_n} \bar{I}(\zeta(v)), \quad c_n = \inf_{\gamma \in \Lambda_n} \max_{v \in U_n} \bar{I}(\gamma(v)).$$

It is obvious that  $c_n \geq b_n$ . For the sake of getting the lower bound of the above minimax values, we give an intersection property which has been proved by Rabinowitz in Lemma 1.44 of [16].

LEMMA 3.1.  $\zeta(D_n) \cap \partial B_\rho \cap E_{n-1}^\perp \neq \emptyset$  for any  $n \in \mathbb{N}$ ,  $\rho < R_n$  and  $\zeta \in \Gamma_n$ .

Next we give the lower bounds for  $b_n$ .

LEMMA 3.2. *There are a positive constant  $C_{17}$  and  $n_0 \in \mathbb{N}$  such that*

$$(3.5) \quad b_n \geq C_{17} n^{2\alpha N - p(N-2)/(N(p-2\alpha))}, \quad n \geq n_0.$$

PROOF. By Lemma 3.1, for any  $\zeta \in \Gamma_n$  and  $\rho < R_n$ , there exists  $v_n \in \zeta(D_n) \cap \partial B_\rho \cap E_{n-1}^\perp$ , then

$$(3.6) \quad \max_{v \in D_n} \bar{I}(\zeta(v)) \geq \bar{I}(v_n) \geq \inf_{v \in \partial B_\rho \cap E_{n-1}^\perp} \bar{I}(v).$$

In view of (g<sub>1</sub>), (g<sub>3</sub>) and (f<sub>3</sub>) in Lemma 2.1, we have

$$(3.7) \quad \int_\Omega |G(x, f(v))| dx \leq C_{18} (\|v\|_{p/\alpha}^{p/\alpha} + 1), \quad v \in E,$$

$$(3.8) \quad \int_\Omega |H(x, f(v))| dx \leq C_{19} (\|v\|_{\sigma/\alpha}^{\sigma/\alpha} + 1), \quad v \in E.$$

In view of (2.7), (3.7) and (3.8),

$$(3.9) \quad \bar{I}(v) \geq \frac{1}{4} \|v\|^2 - C_{20} (\|v\|_{p/\alpha}^{p/\alpha} + 1).$$

By the Gagliardo–Nirenberg inequality, we have

$$(3.10) \quad \|v\|_{p/\alpha} \leq \tau \|v\|^s \|v\|_2^{1-s},$$

where  $\tau$  is a positive constant and  $s = (2p)^{-1}N(p - 2\alpha)$ . If  $v \in E_{n-1}^\perp$ ,

$$(3.11) \quad \|v\|_2^2 \leq \lambda_n^{-1} \|v\|^2.$$

By (3.9), (3.10) and (3.11), if  $v \in \partial B_\rho \cap E_{n-1}^\perp$ ,

$$(3.12) \quad \bar{I}(v) \geq \rho^2 \left( \frac{1}{2} - C_{20} \lambda_n^{(s-1)p/(2\alpha)} \rho^{(p-2\alpha)/\alpha} \right) - C_{20}.$$

In view of (3.12), choose  $\rho_n = (4C_{20})^{\alpha/(2\alpha-p)} \lambda_n^{(1-s)p/(2(p-2\alpha))}$ , then

$$(3.13) \quad \bar{I}(v) \geq \frac{1}{4} \rho_n^2 - C_{20}.$$

It follows from (3.4), (3.6) and (3.13) that (3.5) holds. □

By (H<sub>1</sub>) in Lemma 2.3 and a similar fashion as in the proof of Proposition 10.46 in [17], we have

LEMMA 3.3. *If  $c_n = b_n$  for all  $n \geq n_0$ , where  $n_0$  is a positive integer, there exists a positive constant  $C_{21}$  such that*

$$(3.14) \quad b_n \leq C_{21} n^{2\alpha/(2\alpha-\sigma)}.$$

In view of (h<sub>2</sub>), (3.5) and (3.14), it is impossible that  $c_n = b_n$  for all large  $n$ . Next we can construct critical values of  $\bar{I}$  as follows.

LEMMA 3.4. *Suppose  $c_n > b_n \geq M_2$  for any  $n$  large enough. For any  $\delta \in (0, c_n - b_n)$ , define*

$$(3.15) \quad \Lambda_n(\delta) = \{\gamma \in \Lambda_n : \bar{I}(\gamma(v)) \leq b_n + \delta \text{ for } v \in D_n\}$$

and

$$(3.16) \quad c_n(\delta) = \inf_{\gamma \in \Lambda_n(\delta)} \max_{v \in U_n} \bar{I}(\gamma(v)).$$

Then  $c_n(\delta)$  is a critical value of  $\bar{I}$ .

PROOF. First the definition of  $\Lambda_n(\delta)$  implies that this set is nonempty. By (3.2) and (3.15),

$$(3.17) \quad \Lambda_n(\delta) \subset \Lambda_n, \quad c_n \leq c_n(\delta).$$

By (H<sub>3</sub>) in Lemma 2.3, the Deformation Theorem also holds (see [2]). Suppose  $c_n(\delta)$  is not a critical value of  $\bar{I}$ , choose  $\bar{\varepsilon} := (c_n - b_n - \delta)/2$ , there exists  $\varepsilon \in (0, \bar{\varepsilon})$  and  $\eta \in C([0, 1] \times E, E)$  such that

$$(3.18) \quad \eta(1, v) = v, \quad \bar{I}(v) \notin [c_n(\delta) - \bar{\varepsilon}, c_n(\delta) + \bar{\varepsilon}],$$

$$(3.19) \quad \eta(1, \bar{I}_{c_n(\delta)+\varepsilon}) \subset \bar{I}_{c_n(\delta)-\varepsilon}.$$

By (3.16), there exists  $\gamma \in \Lambda_n(\delta)$  such that

$$(3.20) \quad \max_{v \in U_n} \bar{I}(\gamma(v)) < c_n(\delta) + \varepsilon.$$

Define

$$(3.21) \quad \bar{\gamma}(\cdot) = \eta(1, \gamma(\cdot)).$$

Next we prove  $\bar{\gamma} \in \Lambda_n(\delta)$ . It is obvious that  $\bar{\gamma} \in C(U_n, E)$ . By (3.15) and (3.17),

$$\bar{I}(\gamma(v)) \leq b_n + \delta < c_n - \bar{\varepsilon} \leq c_n(\delta) - \bar{\varepsilon}, \quad v \in D_n,$$

which implies that

$$(3.22) \quad \bar{I}(\gamma(v)) < c_n(\delta) - \bar{\varepsilon}, \quad v \in D_n.$$

In combination with (3.18), (3.21) and (3.22),

$$\bar{\gamma}(v) = \eta(1, \gamma(v)) = \gamma(v), \quad v \in D_n,$$

which yields that

$$(3.23) \quad \bar{\gamma}|_{D_n} \in \Gamma_n \quad \text{and} \quad \bar{I}(\bar{\gamma}(v)) = \bar{I}(\gamma(v)) \leq b_n + \delta, \quad v \in D_n.$$

In view of  $\gamma \in \Lambda_n(\delta)$  and the definitions of  $R_n$  and  $R_{n+1}$ ,

$$(3.24) \quad \gamma(v) = v \quad \text{and} \quad \bar{I}(\gamma(v)) \leq 0, \quad v \in Q_n.$$

Since  $b_n \geq M_2 > 0$  and  $c_n(\delta) \geq c_n > b_n$ , then  $c_n(\delta) > \bar{\varepsilon}$ . It follows from (3.18) and (3.24) that

$$(3.25) \quad \bar{\gamma}(v) = \eta(1, \gamma(v)) = \gamma(v) = v, \quad v \in Q_n.$$

In view of (3.23) and (3.25),  $\bar{\gamma} \in \Lambda_n(\delta)$ . Moreover, by (3.19)–(3.21),

$$\max_{v \in U_n} \bar{I}(\bar{\gamma}(v)) = \max_{v \in U_n} \bar{I}(\eta(1, \gamma(v))) \leq c_n(\delta) - \varepsilon,$$

which is a contradiction to (3.16).  $\square$

**PROOF OF THEOREM 1.1.** Since it is impossible that  $c_n = b_n$  for all large  $n$ , then we can choose a subsequence  $\{n_k\} \subset \mathbb{N}$  such that  $c_{n_k} > b_{n_k}$ . In view of Lemma 3.2,  $c_{n_k} > b_{n_k} > M_2$ , when  $n_k$  is large enough. It follows from (H<sub>2</sub>) in Lemmas 2.3, 3.2 and 3.4 that  $\bar{I}$  has an unbounded sequence of critical values which yields infinitely many solutions for system (1.1).  $\square$

**PROOF OF COROLLARY 1.2.** First, it follows from (g<sub>5</sub>), (h<sub>3</sub>) and (2.1) that  $I$  is an even functional. Arguing as in (H<sub>3</sub>) in Lemma 2.3, we can prove that the functional  $I$  satisfies the  $(C)_c$  condition. Moreover, by a similar fashion as in the proof of Lemma 3.1, there exists a strictly increasing sequence of numbers  $R'_n$  such that  $I(v) \leq 0$  for  $v \in E_n \setminus B_{R'_n}$ . Define

$$\Gamma'_n = \{h \in C(D'_n, E) : h \text{ is odd and } h = \text{id on } \partial B_{R'_n} \cap E_n\},$$

where  $D'_n := \overline{B}_{R'_n} \cap E_n$  and  $b'_n := \inf_{h \in \Gamma'_n} \max_{v \in D'_n} I(h(v))$ . Arguing as in Lemma 3.2, we have  $b'_n \rightarrow \infty$ , as  $n \rightarrow \infty$ . Then there exists  $n_0 \in \mathbb{N}$  such that  $b'_n > 0$ ,  $n \geq n_0$ . If  $n \geq n_0$ , by a standard argument and the Deformation Theorem, we can also prove  $b'_n$  are unbounded critical values of  $I$ .  $\square$

#### 4. Example

In this section, we give one example to illustrate our result.

EXAMPLE 4.1. In system (1.1), let  $\Omega$  be a bounded smooth domain in  $\mathbb{R}^4$  and  $\alpha = 3$ . Let  $g$  is given by (1.5) and  $h(x, t) = t^2$ . Thus all conditions of Theorem 1.1 are satisfied with  $N = 4$ ,  $\kappa = 3$ ,  $\sigma = 7/2$ ,  $p = 13/2$ . By Theorem 1.1, system (1.1) has an unbounded sequence of infinitely many solutions. But the results in [14] cannot be applied to this example.

**Acknowledgements.** The authors would like to thank professors Y.H. Ding, C.G. Liu and W.M. Zou for helpful suggestions and discussions during the *Summer School on Variational Methods and Infinite Dimensional Dynamical System* in Central South University in Changsha.

#### REFERENCES

- [1] S. ADACHI AND T. WATANABE, *Uniqueness of the ground state solutions of quasilinear Schrödinger equations*, *Nonlinear Anal.* **75** (2012), 819–833.
- [2] P. BARTOLO, V. BENCI AND D. FORTUNATO, *Abstract critical point theorems and application to some nonlinear problems with strong resonance at infinity*, *Nonlinear Anal.* **7** (1983), 981–1012.
- [3] R. BARTOLO, A.M. CANDELA AND A. SALVATORE, *Infinitely many radial solutions of a non-homogeneous problem*, *Discrete Contin. Dyn. Syst. Suppl.* (2013), 51–59.
- [4] P. BOLLE, *On the Bolza problem*, *J. Differential Equations* **152** (1999), 274–288.
- [5] P. BOLLE, N. GHOUSSOUB AND H. TEHRANI, *The multiplicity of solutions in nonhomogeneous boundary boundary value problems*, *Manuscripta Math.* **101** (2002), 325–350.
- [6] A.M. CANDELA, G. PALMIERI AND A. SALVATORE, *Radial solutions of semilinear elliptic equations with broken symmetry*, *Topol. Methods Nonlinear Anal.* **27** (2006), 117–132.
- [7] X.D. FANG AND A. SZULKIN, *Multiple solutions for quasilinear Schrödinger equation*, *J. Differential Equations* **254** (2013), 2015–2032.
- [8] S. KURIHARA, *Large-amplitude quasi-solitons in superfluid films*, *J. Phys. Soc. Japan* **50** (1981), 3262–3267.
- [9] E.W. LAEDKE, K.H. SPATSCHEK AND L. STENFLO, *Evolution theorem for a class of perturbed envelope soliton solutions*, *J. Math. Phys.* **24** (1983), 2764–2769.
- [10] J.Q. LIU AND Z.Q. WANG, *Soliton solutions for quasilinear Schrödinger equations I*, *Proc. Amer. Math. Soc.* **131** (2003), 441–448.
- [11] J.Q. LIU, Y. WANG AND Z.Q. WANG, *Solutions for quasilinear Schrödinger equations via the Nehari method*, *Comm. Partial Differential Equations* **29** (2004), 879–892.
- [12] X.Q. LIU, J.Q. LIU AND Z.Q. WANG, *Quasilinear elliptic equations with critical growth via perturbation method*, *J. Differential Equations* **254** (2013), 102–124.
- [13] ———, *Quasilinear elliptic equations via perturbation method*, *Proc. Amer. Math. Soc.* **141** (2013), 253–263.

- [14] X.Q. LIU AND F.K. ZHAO, *Existence of infinitely many solutions for quasilinear elliptic equations perturbed from symmetry*, Adv. Nonlinear Studies **13** (2013), 965–978.
- [15] A. NAKAMURA, *Damping and modification of exciton solitary waves*, J. Phys. Soc. Japan **42** (1977), 1824–1835.
- [16] P. RABINOWITZ, *Multiple critical points of perturbed symmetric functionals*, Trans. Amer. Math. Soc. **272** (1982), 753–769.
- [17] ———, *Minimax Methods in Critical Point Theory with Applications to Differential Equations*, CBMS Reg. Conf. Ser. in Math., vol. 65, Amer. Math. Soc., Providence, RI, 1986.
- [18] D. RUIZ AND G. SICILIANO, *Existence of ground states for a modified nonlinear Schrödinger equation*, Nonlinearity **23** (2010), 1221–1233.
- [19] A. SALVATORE, *Multiple solutions for perturbed elliptic equations in unbounded domains*, Adv. Nonlinear Studies **3** (2003), 1–23.
- [20] M. SCHECHTER AND W. ZOU, *Infinitely many solutions to perturbed elliptic equations*, J. Funct. Anal. **228** (2005), 1–38.
- [21] M. STRUWE, *Infinitely many critical points for functionals which are not even and applications to superlinear boundary value problems*, Manuscripta Math. **32** (1980), 335–364.
- [22] H.T. TEHRANI, *Infinitely many solutions for indefinite semilinear elliptic equations without symmetry*, Comm. Partial Differential Equations **21** (1996), 541–557.
- [23] X. WU, *Multiple solutions for quasilinear Schrödinger equations with a parameter*, J. Differential Equations **256** (2014), 2619–2632.
- [24] X. WU AND K. WU, *Existence of positive solutions, negative solutions and high energy solutions for quasilinear elliptic equations on  $\mathbb{R}^N$* , Nonlinear Anal. RWA **16** (2014), 48–64.
- [25] J. ZHANG, X.H. TANG AND W. ZHANG, *Infinitely many solutions of quasilinear Schrödinger equation with sign-changing potential*, J. Math. Anal. Appl. **420** (2014), 1762–1775.
- [26] ———, *Existence of infinitely many solutions for a quasilinear elliptic equation*, Appl. Math. Lett. **37** (2014), 131–135.

*Manuscript received June 6, 2015*

*accepted July 26, 2015*

LIANG ZHANG  
 School of Mathematical Sciences  
 University of Jinan  
 Jinan, Shandong 250022, P.R. CHINA  
*E-mail address*: mathspaper2012@163.com

XIANHUA TANG  
 School of Mathematics and Statistics  
 Central South University  
 Changsha, Hunan 410083, P.R. CHINA

YI CHEN  
 Department of Mathematics  
 China University of Mining and Technology  
 Xuzhou, 221116, P.R. CHINA