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# STANDING WAVES <br> FOR NONLINEAR SCHRÖDINGER-POISSON EQUATION WITH HIGH FREQUENCY 

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Abstract. We study the existence of ground state and bound state for the following Schrödinger-Poisson equation

$$
\left\{\begin{array}{l}
-\Delta u+V(x) u+\lambda \phi(x) u=\mu u+|u|^{p-1} u, \quad x \in \mathbb{R}^{3},  \tag{P}\\
-\Delta \phi=u^{2}, \quad \lim ^{2} \phi \mid \rightarrow+\infty
\end{array}\right.
$$

where $p \in(3,5), \lambda>0, V \in C\left(\mathbb{R}^{3}, \mathbb{R}^{+}\right)$and $\lim _{|x| \rightarrow+\infty} V(x)=\infty$. By using variational method, we prove that for any $\lambda>0$, there exists $\delta_{1}(\lambda)>0$ such that for $\mu_{1}<\mu<\mu_{1}+\delta_{1}(\lambda)$, problem (P) has a nonnegative ground state with negative energy, which bifurcates from zero solution; problem $(\mathrm{P})$ has a nonnegative bound state with positive energy, which can not bifurcate from zero solution. Here $\mu_{1}$ is the first eigenvalue of $-\Delta+V$. Infinitely many nontrivial bound states are also obtained with the help of a generalized version of symmetric mountain pass theorem.

[^0]
## 1. Introduction

In this paper, we consider the nonlinear Schrödinger-Poisson equation

$$
\left\{\begin{array}{l}
-i \frac{\partial \psi}{\partial t}-\Delta \psi+V(x) \psi+\lambda \phi(x) \psi=|\psi|^{p-1} \psi, \quad(t, x) \in \mathbb{R} \times \mathbb{R}^{3}  \tag{1.1}\\
-\Delta \phi=|\psi|^{2}, \quad \lim _{|x| \rightarrow+\infty} \phi(x)=0
\end{array}\right.
$$

where $\psi: \mathbb{R} \times \mathbb{R}^{3} \rightarrow \mathbb{C}$ denotes the wave function, $\lambda$ is a positive parameter, $p \in(3,5)$, and $V, \phi$ are real valued functions and represent the effective potential and the electric potential, respectively. Problem (1.1) arises from semiconductor theory, see [9], [13] for more physical background.

Recently, there are many papers devoted to looking for standing wave solutions to problem (1.1), that is, $\psi(t, x)=e^{-i \mu t} u(x)$, where $u(x)$ is a real valued function and $\mu \in \mathbb{R}$ denotes the frequency. Then $u(x)$ satisfies the following stationary equation

$$
\left\{\begin{array}{l}
-\Delta u+V(x) u+\lambda \phi(x) u=\mu u+|u|^{p-1} u, \quad x \in \mathbb{R}^{3}  \tag{1.2}\\
-\Delta \phi=u^{2}, \quad \lim _{|x| \rightarrow+\infty} \phi(x)=0
\end{array}\right.
$$

See [2], [6], [7], [10]-[12], [17], [19] and the references therein for all kinds of $V$ and more general nonlinearities.

Throughout this paper, we assume that $V(x)$ satisfies the following condition
(V1) $V \in C\left(\mathbb{R}^{3}, \mathbb{R}^{+}\right)$and $\lim _{|x| \rightarrow+\infty} V(x)=\infty$.
Define

$$
H=\left\{u \in W^{1,2}\left(\mathbb{R}^{3}\right): \int_{\mathbb{R}^{3}} V(x) u^{2} d x<\infty\right\},
$$

with the norm

$$
\|u\|^{2}:=\int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(x) u^{2}\right) d x .
$$

It is known that, under the condition (V1), the embedding $H \hookrightarrow L^{q}\left(\mathbb{R}^{3}\right)(2 \leq$ $q<6)$ is compact. Moreover, there is a sequence of eigenvalues $\left(\mu_{n}\right)$ of $-\Delta+V$ in $H$ such that $0<\mu_{1}<\mu_{2} \leq \mu_{3} \leq \ldots \leq+\infty$ and $H=\operatorname{span}\left\{e_{j}: j \geq 1\right\}$, where $e_{j}$ is the corresponding eigenfunction to $\mu_{j}$ with $\left\|e_{j}\right\|=1$.

For $u \in H$, we denote the unique solution of $-\Delta \phi=u^{2}$ in $D^{1,2}\left(\mathbb{R}^{3}\right)$ by $\phi_{u}$, and

$$
\begin{equation*}
\phi_{u}(x)=\frac{1}{4 \pi} \int_{\mathbb{R}^{3}} \frac{u^{2}(y)}{|x-y|} d y . \tag{1.3}
\end{equation*}
$$

Then equation (1.2) can be rewritten as

$$
\begin{equation*}
-\Delta u+V(x) u+\lambda \phi_{u}(x) u=\mu u+|u|^{p-1} u, \quad x \in \mathbb{R}^{3} . \tag{1.4}
\end{equation*}
$$

Define the energy functional $I_{\mu}: H \rightarrow \mathbb{R}$ by

$$
\begin{align*}
& I_{\mu}(u)=\frac{1}{2} \int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+(V(x)-\mu) u^{2}\right) d x  \tag{1.5}\\
&+\frac{\lambda}{4} \int_{\mathbb{R}^{3}} \phi_{u} u^{2} d x-\frac{1}{p+1} \int_{\mathbb{R}^{3}}|u|^{p+1} d x .
\end{align*}
$$

Then, $I_{\mu} \in C^{1}(H, \mathbb{R})$ and for any $\varphi \in H$ we have that

$$
\begin{align*}
&\left\langle I_{\mu}^{\prime}(u), \varphi\right\rangle=\int_{\mathbb{R}^{3}}(\nabla u \nabla \varphi+(V(x)-\mu) u \varphi) d x  \tag{1.6}\\
&+\lambda \int_{\mathbb{R}^{3}} \phi_{u} u \varphi d x-\int_{\mathbb{R}^{3}}|u|^{p-1} u \varphi d x
\end{align*}
$$

If $u \in H \backslash\{0\}$ and $\left\langle I_{\mu}^{\prime}(u), \varphi\right\rangle=0$ for all $\varphi \in H$, we say that $u$ is a bound state of (1.4). Furthermore, a function $u_{0}$ is called a ground state of (1.4) if $u_{0}$ is a bound state of (1.4) and $I_{\mu}\left(u_{0}\right) \leq I_{\mu}(u)$ for any bound state $u$ of (1.4).

If $\mu<\mu_{1}$ in (1.4), we may define the following equivalent norm on $H$ by

$$
\|u\|_{\mu}^{2}:=\int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+(V(x)-\mu) u^{2}\right) d x .
$$

Then for any $\lambda>0, p \in(3,5)$, and under condition (V1), we can easily prove that (1.4) has a ground state by Nehari manifold method, see [20].

In this paper, we mainly study the existence of bound state, especially ground state, to equation (1.4) for $\mu \geq \mu_{1}$, which is the so called high frequency case in the title. First, we show that (1.4) has a nonnegative bound state with positive energy.

Theorem 1.1. Let $p \in(3,5), \lambda>0$ and (V1) holds. Then there exists $\delta(\lambda)>0$ such that for any $\mu \in\left[\mu_{1}, \mu_{1}+\delta(\lambda)\right.$ ), problem (1.4) has a nonnegative bound state $u_{1, \mu}$ with $I_{\mu}\left(u_{1, \mu}\right)>0$. Moreover, for any sequence $\mu^{(n)} \downarrow \mu_{1}$, there exists $u_{\mu_{1}} \in H$ with $I_{\mu_{1}}^{\prime}\left(u_{\mu_{1}}\right)=0$ and $I_{\mu_{1}}\left(u_{\mu_{1}}\right)>0$, such that $u_{1, \mu^{(n)}} \rightarrow u_{\mu_{1}}$ strongly in $H$.

On basis of Theorem 1.1, we may define the set of all bound states to (1.4):

$$
\mathcal{N}=\left\{u \in H \backslash\{0\}: I_{\mu}^{\prime}(u)=0\right\} \neq \emptyset .
$$

To get the existence of ground state to (1.4), we consider the following minimization problem

$$
\begin{equation*}
c_{0}=\inf \left\{I_{\mu}(u): u \in \mathcal{N}\right\} . \tag{1.7}
\end{equation*}
$$

If we can prove that $c_{0}>-\infty$ and $c_{0} \neq 0$, then by the compactness lemma (see Lemma 2.1) and solving the above minimization problem, the ground state to (1.4) can be obtained. But, it seems not obvious that $c_{0}>-\infty$ and $c_{0} \neq 0$.

Here, let us recall that if $\lambda=0$ in (1.4), by the method of Szulkin and Weth [18] we can prove that $c_{0}>0$. Motivated by Pankov [14], see also [18], we introduce the set

$$
\mathcal{N}_{1}=\left\{u \in H \backslash H_{1}: I_{\mu}^{\prime}(u)=0\right\},
$$

where $H_{1}=\operatorname{span}\left\{e_{1}\right\}$. We first see that $\mathcal{N}=\mathcal{N}_{1}$, that is, $\mathcal{N}_{1}$ contains all bound states of (1.4). Otherwise, if there exists $u \in H_{1} \backslash\{0\}$ such that $I_{\mu}^{\prime}(u)=0$, then noting that $\lambda=0$ and $\mu>\mu_{1}$, we get the following contradiction

$$
0=\left\langle I_{\mu}^{\prime}(u), u\right\rangle=\left(\mu_{1}-\mu\right) \int_{\mathbb{R}^{3}} u^{2} d x-\int_{\mathbb{R}^{3}}|u|^{p+1} d x<0 .
$$

Then we consider the following minimization problem

$$
c_{1}=\inf \left\{I_{\mu}(u): u \in \mathcal{N}_{1}\right\} .
$$

It follows from $\mathcal{N}=\mathcal{N}_{1}$ that $c_{0}=c_{1}$. Next, we want to show $c_{1}>0$. For any fixed $u \in \mathcal{N}_{1}$, let $u=u_{1}+u_{2}, u_{1} \in H_{1}, u_{2} \in H_{2}$, where $H_{2}$ denotes the orthogonal complement of $H_{1}$ in $H$. By the definition of $\mathcal{N}_{1}$, we have $u_{2} \neq 0$. Similar to the proof of Proposition 2.3 in [18], we deduce that for $\lambda=0$,

$$
\begin{equation*}
I_{\mu}(u+w)<I_{\mu}(u), \quad \text { for all } w \in\left\{s u+v: s \geq-1, v \in H_{1}\right\} \backslash\{0\} \tag{1.8}
\end{equation*}
$$

Hence, we get

$$
\begin{equation*}
I_{\mu}(u)=\max \left\{I_{\mu}(v): v \in H_{1} \oplus \mathbb{R}^{+} u_{2}\right\} . \tag{1.9}
\end{equation*}
$$

Since $u_{2} \neq 0$, there exists $s>0$ such that $\left\|s u_{2}\right\|=\rho_{1}$, where $\rho_{1}>0$ is given by Lemma 2.4. Then by (1.9) and Lemma 2.4 we get

$$
I_{\mu}(u) \geq I_{\mu}\left(s u_{2}\right) \geq \alpha_{1}>0 .
$$

Thus, we have $c_{1}>0$, and then $c_{0}>0$ by $c_{0}=c_{1}$.
In case $\lambda>0$ in (1.4), at least for $\lambda>0$ small enough, it can be seen as a small perturbation at $\lambda=0$. So, it seems natural to look for a ground state of (1.4) with positive energy, that is, to get $c_{0}>0$. Following the argument of [18], we find that (1.8) does not hold for $\lambda>0$. This motivates us to doubt that the ground state of (1.4) with positive energy does not exist. In other words, $c_{0}>0$ may do not hold any more. Another main result of the present paper is to prove that (1.4) has a nonnegative ground state with negative energy.

Theorem 1.2. Let $p \in(3,5), \lambda>0$ and (V1) hold. Then there is a $\delta_{1}(\lambda)>0$ and $\delta_{1}(\lambda) \leq \delta(\lambda)$ such that for any $\mu \in\left(\mu_{1}, \mu_{1}+\delta_{1}(\lambda)\right)$, problem (1.4) has a nonnegative ground state $u_{0, \mu}$ with $I_{\mu}\left(u_{0, \mu}\right)<0$. Moreover, for any sequence $\mu^{(n)} \downarrow \mu_{1}, u_{0, \mu^{(n)}} \rightarrow 0$ strongly in $H$.

To prove Theorem 1.1 and Theorem 1.2, by delicate analysis of the nonlocal term in (1.5) we first observe that the energy functional $I_{\mu}$ satisfies mountain pass geometry for $\mu>\mu_{1}$ and near $\mu_{1}$. Then using Ekeland's variational principle
we show that (1.4) has a bound state with negative energy, which also means that $c_{0}<0$. Finally we show that $c_{0}>-\infty$ by the condition (V1).

Next, as a comparison, we recall some known results of (1.4) in the case of $\lambda=0$, i.e.

$$
\begin{equation*}
-\Delta u+V(x) u=\mu u+|u|^{p-1} u, \quad u \in H \tag{1.10}
\end{equation*}
$$

Using the positivity of $e_{1}$, we know that problem (1.10) can not possess any nonnegative bound states for $\mu>\mu_{1}$. Furthermore, by the method of Szulkin and Weth [18] we can prove that problem (1.10) has a ground state with positive energy and the ground state must be sign-changing for $\mu>\mu_{1}$. Theorem 1.1 and Theorem 1.2 show a quite different phenomenon for the problem (1.4) in the case of $\lambda>0$ and $\mu>\mu_{1}$. On the other hand, by using the symmetric property of (1.10) and the condition (V1), problem (1.10) may have infinitely many nontrivial bound states for any $\mu \in \mathbb{R}$ and $1<p<5$. Our next theorem shows that problem (1.4) with $\lambda>0$ also possesses infinitely many nontrivial bound states for any $\mu \in \mathbb{R}$ and $p \in(3,5)$.

Theorem 1.3. Let $p \in(3,5), \lambda>0$ and (V1) hold. Then for any $\mu \in \mathbb{R}$, problem (1.4) has infinitely many nontrivial bound states.

Remark 1.4. We remark here that the condition (V1) was first introduced by Rabinowitz [15], which guarantees the compact embedding of $H \hookrightarrow L^{q}\left(\mathbb{R}^{3}\right)(2 \leq$ $q<6)$. A weaker version has been given by Bartsch and Wang [8]. We emphasize here that all theorems in the present paper will be still true if the condition (V1) is replaced by the weaker condition given by Bartsch and Wang [8].

This paper is organized as follows. In Section 2, we will prove Theorem 1.1. A key step is to prove that the functional $I_{\mu}$ satisfies mountain pass geometry for $\mu$ in a small right neighborhood of $\mu_{1}$. In Section 3, we will study a suitable minimization problem and then use Ekeland variational principle to prove Theorem 1.2. In Section 4, we will use a generalized version of symmetric mountain pass theorem of Rabinowitz [16] to prove Theorem 1.3.

## 2. Nonnegative bound state with positive energy

This section is devoted to the proof of Theorem 1.1. First, we prove a compactness lemma to the functional $I_{\mu}$ on $H$.

Lemma 2.1. Let $p \in(3,5), \lambda, \mu>0$ and (V1) hold. Assume that a sequence $\left(u_{n}\right) \subset H$ satisfies $\left|I_{\mu}\left(u_{n}\right)\right| \leq M<+\infty$ for all $n \in \mathbb{N}$ and $I_{\mu}^{\prime}\left(u_{n}\right) \xrightarrow{n} 0$, then $\left(u_{n}\right)$ has a strongly convergent subsequence in $H$.

Remark 2.2. In view of Lemma 2.1, we say that $I_{\mu}$ satisfies Palais-Smale ((PS) in short) condition.

Proof. By compactness of the embedding $H \hookrightarrow L^{q}\left(\mathbb{R}^{3}\right)(2 \leq q<6)$, it is sufficient to show that $\left(u_{n}\right)$ is bounded in $H$. Choosing $\beta \in(1 /(p+1), 1 / 4)$, then for $n$ large enough we have
(2.1) $M+\left\|u_{n}\right\| \geq I_{\mu}\left(u_{n}\right)-\beta\left\langle I_{\mu}^{\prime}\left(u_{n}\right), u_{n}\right\rangle$

$$
\geq\left(\frac{1}{2}-\beta\right)\left\|u_{n}\right\|^{2}-\left(\frac{1}{2}-\beta\right) \mu \int_{\mathbb{R}^{3}} u_{n}^{2} d x+\left(\beta-\frac{1}{p+1}\right) \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{p+1} d x .
$$

For $\mu>0$, by (V1) there exists $R(\mu)>0$ such that $V(x) \geq 2 \mu$ for all $|x| \geq R(\mu)$.
Then

$$
\begin{equation*}
\int_{|x| \geq R(\mu)} u_{n}^{2} d x \leq \int_{|x| \geq R(\mu)} \frac{V(x)}{2 \mu} u_{n}^{2} d x \leq \frac{1}{2 \mu}\left\|u_{n}\right\|^{2} . \tag{2.2}
\end{equation*}
$$

On the other hand, by Young's inequality we get for any $\varepsilon>0$,

$$
\begin{align*}
\int_{|x| \leq R(\mu)} u_{n}^{2} d x & \leq C(\mu)\left(\int_{|x| \leq R(\mu)}\left|u_{n}\right|^{p+1} d x\right)^{2 /(p+1)}  \tag{2.3}\\
& \leq \varepsilon \int_{\mathbb{R}^{3}}\left|u_{n}\right|^{p+1} d x+C(\varepsilon, \mu)
\end{align*}
$$

Taking $\varepsilon=(\beta-1 /(p+1)) /((1 / 2-\beta) \mu)$ in (2.3) and combining (2.1)-(2.3) we deduce that

$$
\begin{equation*}
M+\left\|u_{n}\right\| \geq \frac{1}{2}\left(\frac{1}{2}-\beta\right)\left\|u_{n}\right\|^{2}-C(\beta, \mu) \tag{2.4}
\end{equation*}
$$

Thus, $\left(u_{n}\right)$ is bounded in $H$.
REmARK 2.3. The proof of (2.4) will play a crucial role to show that $c_{0}>$ $-\infty$, where $c_{0}$ is defined by (1.7), see the proof of Theorem 1.2 , step 2, in Section 3 .

In order to prove Theorem 1.1, we use the following classical mountain pass Lemma due to Ambrosetti and Rabinowitz [3].

Lemma 2.4. Let $E$ be a real Banach space and the functional $I \in C^{1}(E, \mathbb{R})$. Suppose that $I(0)=0$ and
(a) there are constants $\rho, \alpha>0$ such that $\left.I\right|_{\partial B_{\rho}} \geq \alpha$; and
(b) there is $\bar{u} \in E \backslash \overline{B_{\rho}}$ such that $I(\bar{u})<0$.

Let $c$ be defined by

$$
c=\inf _{g \in \Gamma} \max _{u \in g[0,1]} I(u) \quad \text { with } \Gamma=\{g \in C([0,1], E): g(0)=0, g(1)=\bar{u}\} .
$$

If I satisfies the (PS) condition, then I possesses a critical value $c \geq \alpha$.
In view of Lemma 2.1, a key step of applying Lemma 2.4 is to verify that the functional $I_{\mu}$ satisfy mountain pass geometry, i.e. (a) and (b) in Lemma 2.4. In the case of $0<\mu<\mu_{1}$, it is easy to prove that the functional $I_{\mu}$ defined on $H$
satisfies mountain pass geometry. However, for $\mu>\mu_{1}$, it is a difficult issue to prove that the mountain pass geometry holds for $I_{\mu}$. We have to analyze the structure of the functional $I_{\mu}$ delicately and find that the competing of the Poisson term $\int_{\mathbb{R}^{3}} \phi_{u} u^{2} d x$ and the nonlinear term $\int_{\mathbb{R}^{3}}|u|^{p+1} d x$ may produce a new phenomenon to the geometric structure of the functional $I_{\mu}$. And we manage to get the mountain pass geometry of the functional $I_{\mu}$ for $\mu$ in a small right neighbourhood of $\mu_{1}$. The more precise statement is the following lemma.

Lemma 2.5. Let $p \in(3,5), \lambda>0$ and (V1) hold. Then we have the following conclusions:
(a) If $0<\mu<\mu_{1}$, then 0 is a local minimum of $I_{\mu}$.
(b) There are positive constants $\delta(\lambda), \rho(\lambda)$ and $\alpha(\lambda)$ such that, for any $\mu \in$ $\left[\mu_{1}, \mu_{1}+\delta(\lambda)\right),\left.I_{\mu}\right|_{\partial B_{\rho(\lambda)}} \geq \alpha(\lambda)$.
(c) There is $\bar{u} \in H^{1}\left(\mathbb{R}^{3}\right)$ with $\|\bar{u}\|>\rho(\lambda)$ such that $I_{\mu}(\bar{u})<0$ for any $\mu>0$.

Proof. (a) For any $u \in H \backslash\{0\}$, from $p \in(3,5), \lambda>0,0<\mu<\mu_{1}$ and the continuity of the Sobolev embedding of $H$ in $L^{p+1}\left(\mathbb{R}^{3}\right)$, we deduce that

$$
\begin{aligned}
I_{\mu}(u) & =\frac{1}{2}\|u\|^{2}+\frac{\lambda}{4} \int_{\mathbb{R}^{3}} \phi_{u} u^{2} d x-\frac{1}{p+1} \int_{\mathbb{R}^{3}}|u|^{p+1} d x-\frac{\mu}{2} \int_{\mathbb{R}^{3}} u^{2} d x \\
& \geq \frac{1}{2}\|u\|^{2}-C\|u\|^{p+1}-\frac{\mu}{2 \mu_{1}}\|u\|^{2}=\|u\|^{2}\left(\frac{1}{2}-\frac{\mu}{2 \mu_{1}}-C\|u\|^{p-1}\right) .
\end{aligned}
$$

Choosing $\rho_{0}=\|u\|$ small enough such that $C \rho_{0}^{p-1} \leq\left(1-\mu / \mu_{1}\right) / 4$, we obtain that

$$
I_{\mu}(u) \geq \frac{1}{4}\left(1-\frac{\mu}{\mu_{1}}\right) \rho_{0}^{2}
$$

Therefore the conclusion (a) follows.
(b) Our goal is to prove that for any $\lambda>0$, there exist $\rho(\lambda), \alpha(\lambda), \delta(\lambda)>0$ such that for any $\mu_{1} \leq \mu<\mu_{1}+\delta(\lambda)$,

$$
\begin{equation*}
I_{\mu}(u) \geq \alpha(\lambda), \quad \text { for all } u \in H \text { with }\|u\|=\rho(\lambda) \tag{2.5}
\end{equation*}
$$

We first show that for any $\lambda>0$, there exist constants $\rho(\lambda), \alpha(\lambda)>0$ such that

$$
\begin{equation*}
I_{\mu_{1}}(u) \geq \alpha(\lambda), \quad \text { for all } u \in H \text { with }\|u\|=\rho(\lambda) . \tag{2.6}
\end{equation*}
$$

Define

$$
F(u)=\int_{\mathbb{R}^{3}} \phi_{u} u^{2} d x \quad \text { for } u \in H
$$

Then for any $u \in H$, there exists $t=t(u) \in \mathbb{R}$ and $v \in H_{2}$ such that $u=t e_{1}+v$ and $\|u\|^{2}=t^{2}+\|v\|^{2}$.

By the mean value theorem we have

$$
\begin{align*}
\left|F(u)-F\left(t e_{1}\right)\right| & \leq\left|\left\langle F^{\prime}\left(t e_{1}+\theta v\right), v\right\rangle\right|, \quad \text { for } \theta \in[0,1]  \tag{2.7}\\
& =4\left|\int_{\mathbb{R}^{3}} \phi_{t e_{1}+\theta v}\left(t e_{1}+\theta v\right) v d x\right| \\
& \leq 4\left\|\phi_{t e_{1}+\theta v}\right\|_{L^{6}}\left\|t e_{1}+\theta v\right\|_{L^{3}}\|v\|_{L^{2}} \\
& \leq C\left\|t e_{1}+\theta v\right\|^{3}\|v\| \leq C\left(|t|^{3}\|v\|+\|v\|^{4}\right) .
\end{align*}
$$

Therefore,

$$
\begin{align*}
I_{\mu_{1}}(u)= & \frac{1}{2}\|u\|^{2}-\frac{\mu_{1}}{2} \int_{\mathbb{R}^{3}} u^{2} d x  \tag{2.8}\\
& +\frac{\lambda}{4}\left(F(u)-F\left(t e_{1}\right)+F\left(t e_{1}\right)\right)-\frac{1}{p+1} \int_{\mathbb{R}^{3}}|u|^{p+1} d x \\
\geq & \frac{1}{2}\left(t^{2}+\|v\|^{2}\right)-\frac{1}{2}\left(t^{2}+\frac{\mu_{1}}{\mu_{2}}\|v\|^{2}\right)-C_{1} \lambda\left(|t|^{3}\|v\|+\|v\|^{4}\right) \\
& +\frac{\lambda}{4} t^{4} \int_{\mathbb{R}^{3}} \phi_{e_{1}} e_{1}^{2} d x-C \int_{\mathbb{R}^{3}}\left(\left|t e_{1}\right|^{p+1}+|v|^{p+1}\right) d x \\
= & \frac{1}{2}\left(1-\frac{\mu_{1}}{\mu_{2}}\right)\|v\|^{2}-C_{1} \lambda\left(|t|^{3}\|v\|+\|v\|^{4}\right) \\
& +C_{2} \lambda t^{4}-C_{3}|t|^{p+1}-C_{4}\|v\|^{p+1} .
\end{align*}
$$

Setting $\varepsilon_{1}=\left(1-\mu_{1} / \mu_{2}\right) /\left(4 C_{1} \lambda\right)$, then

$$
\begin{align*}
\frac{1}{2}\left(1-\frac{\mu_{1}}{\mu_{2}}\right) & \|v\|^{2}-C_{1} \lambda|t|^{3}\|v\|  \tag{2.9}\\
& \geq \frac{1}{2}\left(1-\frac{\mu_{1}}{\mu_{2}}\right)\|v\|^{2}-C_{1} \lambda\left(\varepsilon_{1}\|v\|^{2}+\frac{1}{\varepsilon_{1}} t^{6}\right) \\
& =\frac{1}{4}\left(1-\frac{\mu_{1}}{\mu_{2}}\right)\|v\|^{2}-C_{5} \lambda^{2} t^{6} .
\end{align*}
$$

Noting that $3<p<5$, for

$$
|t| \leq \min \left\{1,\left(\frac{C_{2} \lambda}{2\left(C_{3}+C_{5} \lambda^{2}\right)}\right)^{1 /(p-3)}\right\} \triangleq \xi
$$

and

$$
\|v\| \leq \min \left\{1,\left(\frac{\left(1-\mu_{1} / \mu_{2}\right) / 8}{C_{1} \lambda+C_{4}}\right)^{1 / 2}\right\} \triangleq \eta
$$

we have
(2.10) $-C_{5} \lambda^{2} t^{6}+C_{2} \lambda t^{4}-C_{3}|t|^{p+1} \geq C_{2} \lambda t^{4}-\left(C_{3}+C_{5} \lambda^{2}\right)|t|^{p+1} \geq \frac{C_{2}}{2} \lambda t^{4}$,
and

$$
\begin{align*}
\frac{1}{4}\left(1-\frac{\mu_{1}}{\mu_{2}}\right) & \|v\|^{2}-C_{1} \lambda\|v\|^{4}-C_{4}\|v\|^{p+1}  \tag{2.11}\\
& \geq \frac{1}{4}\left(1-\frac{\mu_{1}}{\mu_{2}}\right)\|v\|^{2}-C_{1} \lambda\|v\|^{4}-C_{4}\|v\|^{4} \\
& \geq \frac{1}{8}\left(1-\frac{\mu_{1}}{\mu_{2}}\right)\|v\|^{2} \geq \frac{1}{8}\left(1-\frac{\mu_{1}}{\mu_{2}}\right)\|v\|^{4}
\end{align*}
$$

Combining (2.8)-(2.11) we get

$$
\begin{equation*}
I_{\mu_{1}}(u) \geq \frac{C_{2}}{2} \lambda t^{4}+\frac{1}{8}\left(1-\frac{\mu_{1}}{\mu_{2}}\right)\|v\|^{4} \geq C(\lambda)\left(t^{2}+\|v\|^{2}\right)^{2}=C(\lambda)\|u\|^{4} \tag{2.12}
\end{equation*}
$$

Let $\rho(\lambda)=\min \{\xi, \eta\}$. Then by (2.12) we see that (2.6) holds.
For $\mu>\mu_{1}$, we have

$$
I_{\mu}(u)=I_{\mu_{1}}(u)-\frac{1}{2}\left(\mu-\mu_{1}\right) \int_{\mathbb{R}^{3}} u^{2} d x \geq I_{\mu_{1}}(u)-\left(\mu-\mu_{1}\right) C\|u\|^{2} .
$$

This and (2.6) imply that there is $\delta(\lambda)>0$ such that (2.5) holds.
(c) Let $e_{0} \in H \backslash\{0\}$ and $s>0$. Then we have that

$$
\begin{aligned}
I_{\mu}\left(s e_{0}\right) & =\frac{s^{2}}{2}\left(\left\|e_{0}\right\|^{2}-\mu \int_{\mathbb{R}^{3}} e_{0}^{2} d x\right)+\frac{\lambda s^{4}}{4} \int_{\mathbb{R}^{3}} \phi_{e_{0}} e_{0}^{2} d x-\frac{s^{p+1}}{p+1} \int_{\mathbb{R}^{3}}\left|e_{0}\right|^{p+1} d x \\
& \leq \frac{s^{2}}{2}\left\|e_{0}\right\|^{2}+\frac{\lambda s^{4}}{4} \int_{\mathbb{R}^{3}} \phi_{e_{0}} e_{0}^{2} d x-\frac{s^{p+1}}{p+1} \int_{\mathbb{R}^{3}}\left|e_{0}\right|^{p+1} d x
\end{aligned}
$$

Since $\left\|e_{0}\right\|^{2}, \int_{\mathbb{R}^{3}} \phi_{e_{0}} e_{0}^{2} d x$ and $\int_{\mathbb{R}^{3}}\left|e_{0}\right|^{p+1} d x$ are fixed and positive, the fact of $p \in(3,5)$ implies that there exists $s_{0}>0$ such that

$$
\left\|s_{0} e_{0}\right\|>\rho(\lambda) \quad \text { and } \quad I_{\mu}\left(s_{0} e_{0}\right)<0
$$

The conclusion (c) follows from choosing $\bar{u}=s_{0} e_{0}$.
Proof of Theorem 1.1. We denote

$$
c_{1, \mu}=\inf _{\gamma \in \Gamma} \max _{t \in[0,1]} I_{\mu}(\gamma(t)) \quad \text { with } \Gamma=\left\{\gamma \in C([0,1], H): \gamma(0)=0, \gamma(1)=s_{0} e_{0}\right\}
$$

By Lemmas 2.1 and 2.5, the mountain pass Lemma 2.4 implies that $c_{1, \mu}$ is a critical value of $I_{\mu}$ and $c_{1, \mu}>0$. The proof of nonnegativity for at least one of the corresponding critical point is inspired by the idea of [1]. In fact, since $I_{\mu}(u)=I_{\mu}(|u|)$ in $H$, for every $n \in \mathbb{N}$, there exists $\gamma_{n} \in \Gamma$ with $\gamma_{n}(t) \geq 0$ (almost everywhere in $\mathbb{R}^{3}$ ) for all $t \in[0,1]$ such that

$$
\begin{equation*}
c_{1, \mu} \leq \max _{t \in[0,1]} I_{\mu}\left(\gamma_{n}(t)\right)<c_{1, \mu}+\frac{1}{n} \tag{2.13}
\end{equation*}
$$

Consequently, by means of Ekeland's principle (see for instance [5]), there exists $\gamma_{n}^{*} \in \Gamma$ with the following properties:

$$
\left\{\begin{array}{l}
c_{1, \mu} \leq \max _{t \in[0,1]} I_{\mu}\left(\gamma_{n}^{*}(t)\right) \leq \max _{t \in[0,1]} I_{\mu}\left(\gamma_{n}(t)\right)<c_{1, \mu}+\frac{1}{n}  \tag{2.14}\\
\left.\left.\max _{t \in[0,1]} \| \gamma_{n}(t)\right)-\gamma_{n}^{*}(t)\right) \|<\frac{1}{\sqrt{n}} \\
\text { there exists } t_{n} \in[0,1] \text { such that } z_{n}=\gamma_{n}^{*}\left(t_{n}\right) \text { satisfies: } \\
I_{\mu}\left(z_{n}\right)=\max _{t \in[0,1]} I_{\mu}\left(\gamma_{n}^{*}(t)\right), \text { and }\left\|I_{\mu}^{\prime}\left(z_{n}\right)\right\| \leq \frac{1}{\sqrt{n}}
\end{array}\right.
$$

By Lemma 2.1 we get a convergent subsequence (still denoted by $\left(z_{n}\right)_{n \in \mathbb{N}}$ ). We may assume that $z_{n} \rightarrow z$ in $H^{1}\left(\mathbb{R}^{3}\right)$ as $n \rightarrow \infty$. On the other hand, by (2.14), we also arrive at $\gamma_{n}\left(t_{n}\right) \rightarrow z$ in $H^{1}\left(\mathbb{R}^{3}\right)$ as $n \rightarrow \infty$. Since $\gamma_{n}(t) \geq 0$, we conclude that $z \geq 0, z \not \equiv 0$ in $\mathbb{R}^{3}$ with $I_{\mu}(z)>0$ and it is a nonnegative bound state of problem (1.4).

Let $u_{1, \mu}$ be the nonnegative bound state given by the above proof, that is,

$$
I_{\mu}^{\prime}\left(u_{1, \mu}\right)=0 \quad \text { and } \quad I_{\mu}\left(u_{1, \mu}\right)=c_{1, \mu} .
$$

We claim that for any sequence $\mu^{(n)} \downarrow \mu_{1}$, there exists $u_{\mu_{1}} \in H$ with $I_{\mu_{1}}^{\prime}\left(u_{\mu_{1}}\right)=0$ and $I_{\mu_{1}}\left(u_{\mu_{1}}\right)>0$, such that $u_{1, \mu^{(n)}} \rightarrow u_{\mu_{1}}$ strongly in $H$. In fact, by the definition of $c_{1, \mu}$ and the proof of Lemma 2.5(c), we deduce that for $n$ large enough,

$$
\begin{aligned}
0<\alpha(\lambda) & \leq c_{1, \mu^{(n)}} \leq \max _{s \geq 0} I_{\mu^{(n)}}\left(s e_{0}\right) \\
& \leq \max _{s \geq 0}\left\{\frac{s^{2}}{2}\left\|e_{0}\right\|^{2}+\frac{\lambda s^{4}}{4} \int_{\mathbb{R}^{3}} \phi_{e_{0}} e_{0}^{2} d x-\frac{s^{p+1}}{p+1} \int_{\mathbb{R}^{3}}\left|e_{0}\right|^{p+1} d x\right\} .
\end{aligned}
$$

Thus, the critical value sequence $\left\{c_{1, \mu^{(n)}}\right\}$ is bounded from above and below. Then by the proof of Lemma 2.1 we see that the claim holds.

## 3. Nonnegative ground state with negative energy

In this section, we prove another main result of this paper, that is, problem (1.4) has a nonnegative ground state with negative energy.

Proof of Theorem 1.2. We divide the proof into three steps.
Step 1. We claim that there exists $w \in H$ such that

$$
I_{\mu}^{\prime}(w)=0 \quad \text { and } \quad I_{\mu}(w)<0
$$

For $\rho(\lambda)>0$ given in the proof of Lemma 2.5, we set

$$
\begin{equation*}
c_{2, \mu}=\inf \left\{I_{\mu}(u):\|u\| \leq \rho(\lambda)\right\} \tag{3.1}
\end{equation*}
$$

It is clear that $c_{2, \mu}>-\infty$.

We will show that $c_{2, \mu}<0$. In fact, since $\mu>\mu_{1}$ we have for $t>0$ small enough,

$$
\begin{align*}
I_{\mu}\left(t e_{1}\right)= & \frac{t^{2}}{2}\left(\left\|e_{1}\right\|^{2}-\mu \int_{\mathbb{R}^{3}} e_{1}^{2} d x\right)  \tag{3.2}\\
& +\frac{\lambda}{4} t^{4} \int_{\mathbb{R}^{3}} \phi_{e_{1}} e_{1}^{2} d x-\frac{1}{p+1} t^{p+1} \int_{\mathbb{R}^{3}}\left|e_{1}\right|^{p+1} d x \\
= & \frac{t^{2}}{2}\left(1-\frac{\mu}{\mu_{1}}\right)+C \lambda t^{4}-C t^{p+1}<0 .
\end{align*}
$$

This implies that $c_{2, \mu}<0$. Note that if $\left(u_{n}\right)$ is a minimizing sequence of $c_{2, \mu}$, then $\left(\left|u_{n}\right|\right)$ is also a minimizing of $c_{2, \mu}$. Therefore by (3.1) and Ekeland's variational principle there exists a sequence $\left(u_{n}\right)$ in $H$ and $u_{n} \geq 0$ such that

$$
I_{\mu}\left(u_{n}\right) \xrightarrow{n} c_{2, \mu} \quad \text { and } \quad I_{\mu}^{\prime}\left(u_{n}\right) \xrightarrow{n} 0 .
$$

It follows from Lemma 2.1 that there exists $w \in H$ and $w \geq 0$ such that $I_{\mu}^{\prime}(w)=0$ and $I_{\mu}(w)<0$.

Step 2. We claim that there exists $u_{0, \mu} \in H$ such that

$$
I_{\mu}^{\prime}\left(u_{0, \mu}\right)=0 \quad \text { and } \quad I_{\mu}\left(u_{0, \mu}\right)=c_{0, \mu}<0
$$

where

$$
c_{0, \mu}=\inf \left\{I_{\mu}(u): u \in \mathcal{N}\right\}, \quad \mathcal{N}=\left\{u \in H \backslash\{0\}: I_{\mu}^{\prime}(u)=0\right\} .
$$

By Step 1 we know that $\mathcal{N} \neq \emptyset$ and $c_{0, \mu}<0$. We will show that $c_{0, \mu}>-\infty$. For any $u \in \mathcal{N}$, similar to the proof of (2.4), we have that

$$
I_{\mu}(u)=I_{\mu}(u)-\beta\left\langle I_{\mu}^{\prime}(u), u\right\rangle \geq \frac{1}{2}\left(\frac{1}{2}-\beta\right)\|u\|^{2}-C(\beta, \mu)
$$

where $\beta \in(1 /(p+1), 1 / 4)$. Thus $c_{0, \mu}>-\infty$. Choosing $u_{n} \in \mathcal{N}$ such that $I_{\mu}\left(u_{n}\right) \xrightarrow{n} c_{0, \mu}$ and $I_{\mu}^{\prime}\left(u_{n}\right) \xrightarrow{n} 0$, then by Lemma 2.1 there exists $u_{0, \mu} \in H$ such that $I_{\mu}^{\prime}\left(u_{0, \mu}\right)=0$ and $I_{\mu}\left(u_{0, \mu}\right)=c_{0, \mu}$, which means that problem (1.4) has a ground state $u_{0, \mu}$ with $I_{\mu}\left(u_{0, \mu}\right)<0$.

Step 3. We claim that the solutions given by Steps 1 and 2 coincide. The proof is divided into two steps. In the first place, for any $u \neq 0$ and $u$ is a solution of (1.4) with $\mu=\mu_{1}$, we have that

$$
\|u\|^{2}-\mu_{1} \int_{\mathbb{R}^{3}}|u|^{2} d x+\int_{\mathbb{R}^{3}} \phi_{u} u^{2} d x=\int_{\mathbb{R}^{3}}|u|^{p+1} d x
$$

and hence

$$
I_{\mu_{1}}(u)=\left(\frac{1}{2}-\frac{1}{p+1}\right)\left(\|u\|^{2}-\mu_{1} \int_{\mathbb{R}^{3}}|u|^{2} d x\right)+\left(\frac{1}{4}-\frac{1}{p+1}\right) \int_{\mathbb{R}^{3}} \phi_{u} u^{2} d x
$$

Since $\|u\|^{2} \geq \mu_{1} \int_{\mathbb{R}^{3}}|u|^{2} d x$ for any $u \in H$,

$$
I_{\mu_{1}}(u) \geq\left(\frac{1}{4}-\frac{1}{p+1}\right) \int_{\mathbb{R}^{3}} \phi_{u} u^{2} d x>0
$$

In the second place, let $u_{0, \mu}$ be the ground state given by the above Step 2. For any sequence $\mu^{(n)} \downarrow \mu_{1}$, by the proof of Lemma 2.1 we deduce that $u_{0, \mu^{(n)}}$ must converge to a critical point of $I_{\mu_{1}}$ with non-positive energy, that is, 0 . Hence there is $\delta_{1}(\lambda)>0$ and $\delta_{1}(\lambda) \leq \delta(\lambda)$ such that for $\mu \in\left(\mu_{1}, \mu_{1}+\delta_{1}(\lambda)\right),\left\|u_{0, \mu}\right\|<\rho(\lambda)$, which implies that $c_{0, \mu}=c_{2, \mu}$. Hence we can conclude that $w=u_{0, \mu}$, which is a nonnegative ground state of (1.4) for $\mu \in\left(\mu_{1}, \mu_{1}+\delta_{1}(\lambda)\right)$. The proof of Theorem 1.2 is complete.

Remark 3.1. (a) Combining Theorems 1.1 and 1.2, we know that (1.4) has at least two nonnegative bound states $u_{0, \mu}$ and $u_{1, \mu}$ for $\mu \in\left(\mu_{1}, \mu_{1}+\delta_{1}(\lambda)\right)$. Theorem 1.2 implies that the ground state $u_{0, \mu}$ bifurcates from zero solution. But Theorem 1.1 implies that the bound state $u_{1, \mu}$ can not bifurcate from zero solution.
(b) In the proof of Step 2 of Theorem 1.2, we do not use the Nehari manifold method, that is, we do not consider the following minimization problem

$$
c_{0}=\inf \left\{I_{\mu}(u): u \in \mathcal{M}\right\}, \quad \mathcal{M}=\left\{u \in H \backslash\{0\}:\left\langle I_{\mu}^{\prime}(u), u\right\rangle=0\right\} .
$$

This is because for $\mu>\mu_{1}$, we cannot deduce that $0 \notin \partial \mathcal{M}$.

## 4. Infinitely many nontrivial bound states

In this section, we will prove Theorem 1.3 by using the condition (V1) and the fact that the problem (1.4) is symmetric with respect to $u \in H$. We start with the following Theorem 4.1 from Rabinowitz [16].

Theorem 4.1. Let $E$ be an infinite dimensional Banach space and let $I \in$ $C^{1}(E, \mathbb{R})$ be even and $I$ satisfies $(\mathrm{PS})$ conditions. Suppose that $I(0)=0$ and $E=Y \oplus X$, where $Y$ is finite dimensional and I satisfies:
(a) there are constants $\rho, \alpha>0$ such that $\left.I\right|_{\partial B_{\rho} \cap X} \geq \alpha$ and
(b) for each finite dimensional subspace $\widetilde{E} \subset E$, there is an $R=R(\widetilde{E})$ such that $I \leq 0$ on $\widetilde{E} \backslash B_{R(\widetilde{E})}$.
Then I possesses an unbounded sequence of critical values.
We are now in a position to use this theorem to prove that the problem (1.4) has infinitely many nontrivial bound states.

Proof of Theorem 1.3. If $\mu<\mu_{1}$, then we may use standard symmetric mountain pass theorem [3], [16] to get the conclusion. In the following, we may assume without loss of generality that $\mu_{k} \leq \mu<\mu_{k+1}$, where $\mu_{k}$ is the $k$-th eigenvalue of $-\Delta+V$ in $H$. From Lemma 2.1, we know that $I_{\mu}$ satisfies (PS) condition. Clearly $I_{\mu}(0)=0$. Choosing $E=H, Y=\operatorname{span}\left\{e_{1}, \ldots, e_{k}\right\}$ and $X=Y^{\perp}$, we are in a position to verify (a) and (b) of Theorem 4.1.
( $\mathrm{a}^{\prime}$ ) For any $u \in X$, since $\mu_{k} \leq \mu<\mu_{k+1}$, we have that

$$
\int_{\mathbb{R}^{3}}\left(|\nabla u|^{2}+V(x) u^{2}\right) d x \geq \mu_{k+1} \int_{\mathbb{R}^{3}} u^{2} d x .
$$

Therefore

$$
\begin{aligned}
I_{\mu}(u) & \geq \frac{1}{2}\left(1-\frac{\mu}{\mu_{k+1}}\right)\|u\|^{2}+\frac{\lambda}{4} \int \phi_{u} u^{2} d x-\frac{C_{p+1}}{p+1}\|u\|^{p+1} \\
& \geq \frac{1}{2}\left(1-\frac{\mu}{\mu_{k+1}}\right)\|u\|^{2}\left(1-\frac{2 \mu_{k+1} C_{p+1}}{(p+1)\left(\mu_{k+1}-\mu\right)}\|u\|^{p-1}\right) .
\end{aligned}
$$

Hence there is a $\rho>0$ such that for $\|u\|=\rho$,

$$
I_{\mu}(u) \geq \frac{1}{4}\left(1-\frac{\mu}{\mu_{k+1}}\right) \rho^{2} .
$$

(b') For each finite dimensional $\widetilde{E} \subset H$ and for any $v \in \widetilde{E}$,

$$
\begin{aligned}
I_{\mu}(v) & =\frac{1}{2}\|v\|^{2}-\frac{\mu}{2} \int v^{2} d x+\frac{\lambda}{4} \int \phi_{v} v^{2} d x-\frac{1}{p+1} \int|v|^{p+1} d x \\
& \leq \frac{1}{2}\|v\|^{2}-\frac{\mu}{2} \int v^{2} d x+\frac{\lambda C}{4}\|v\|^{4}-\frac{1}{p+1} \int|v|^{p+1} d x .
\end{aligned}
$$

Since $\widetilde{E}$ is finite dimensional and $p \in(3,5)$, we see that there is $R:=R(\widetilde{E})$ such that for all $v \in \widetilde{E} \backslash B_{R(\widetilde{E})}, I_{\mu}(v) \leq 0$.

Now using Theorem 4.1 we know that $I_{\mu}$ possesses an unbounded sequence of critical values and hence the problem (1.4) has infinitely many nontrivial bound states.

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## References

[1] S. Alama and G. Tarantello, On semilinear elliptic equations with indefinite nonlinearities, Calc. Var. Partial Differential Equations 1 (1993), 439-475.
[2] A. Ambrosetti, On Schrödinger-Poisson systems, Milan J. Math. 76 (2008), 257-274.
[3] A. Ambrosetti and P.H. Rabinowitz, Dual variational methods in critical point theory and applications, J. Funct. Anal. 14 (1973), 349-381.
[4] A. Ambrosetti and D. Ruiz, Multiple bound states for the Schrödinger-Poisson problem, Commun. Contemp. Math. 10 (2008), 391-404.
[5] T. Aubin and I. Ekeland, Applied Nonlinear Analysis, Wiley, New York, 1984.
[6] A. Azzollini, P. D'Avenia and A. Pomponio, On the Schrödinger-Maxwell equations under the effect of a general nonlinear term, Ann. Inst. H. Poincaré Anal. Non Linéaire 27 (2010), 779-791.
[7] A. Azzollini and A. Pomponio, Ground state solutions for the nonlinear SchrödingerMaxwell equations, J. Math. Anal. Appl. 345 (2008), 90-108.
[8] T. Bartsch and Z.-Q. Wang, Existence and multiplicity results for some superlinear elliptic problems on $\mathbb{R}^{N}$, Comm. Partial Differential Equations 20 (1995), 1725-1741.
[9] V. Benci and D. Fortunato, An eigenvalue problem for the Schrödinger-Maxwell equations, Topol. Methods Nonlinear Anal. 11 (1998), 283-293.
[10] T. D'Aprile and J. Wei, Standing waves in the Maxwell-Schrödinger equation and an optimal configuration problem, Calc. Var. Partial Differential Equations 25 (2006), 105137.
[11] L. Huang, E.M. Rocha and J. Chen, Two positive solutions of a class of SchrödingerPoisson system with indefinite nonlinearity, J. Differential Equations 255 (2013), 24632483.
[12] G.B. Li, S.J. Peng and S.S. Yan, Infinitely many positive solutions for the nonlinear Schrödinger-Poisson system, Commun. Contemp. Math. 12 (2010), 1069-1092.
[13] P.A. Markowich, C. Ringhofer and C. Schmeiser, Semiconductor Equations, Springer-Verlag, New York, 1990.
[14] A. Pankov, Periodic nonlinear Schrödinger equation with application to photonic crystals, Milan J. Math. 73 (2005), 259-287.
[15] P. Rabinowitz, On a class of nonlinear Schrödinger equations, Z. Angew. Math. Phys. 43 (1992), 270-291.
[16] _, Minimax Methods in Critical Point Theory with Applications to Differential Equations, CBMS 65 (American Mathematical Society), Providence, 1986.
[17] D. Ruiz, The Schrödinger-Poissom equation under the effect of a nonlinear local term, J. Funct. Anal. 237 (2006), 655-674.
[18] A. Szulkin and T. Weth, Ground state solutions for some indefinite variational problems, J. Funct. Anal. 257 (2009), 3802-3822.
[19] Z.P. Wang and H.S. Zhou, Positive solution for a nonlinear stationary SchrödingerPoisson system in $\mathbb{R}^{3}$, Discrete Contin. Dyn. Syst. 18 (2007), 809-816.
[20] M. Willem, Minimax Theorems, Birkhäuser, Boston, 1996.

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