# DIMENSION OF ATTRACTORS AND INVARIANT SETS IN REACTION DIFFUSION EQUATIONS 

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#### Abstract

Under fairly general assumptions, we prove that every compact invariant set $\mathcal{I}$ of the semiflow generated by the semilinear reaction diffusion


 equation$$
\begin{aligned}
u_{t}+\beta(x) u-\Delta u & =f(x, u), & & (t, x) \in[0,+\infty[\times \Omega, \\
u & =0, & & (t, x) \in[0,+\infty \times \partial \Omega
\end{aligned}
$$

in $H_{0}^{1}(\Omega)$ has finite Hausdorff dimension. Here $\Omega$ is an arbitrary, possibly unbounded, domain in $\mathbb{R}^{3}$ and $f(x, u)$ is a nonlinearity of subcritical growth. The nonlinearity $f(x, u)$ needs not to satisfy any dissipativeness assumption and the invariant subset $\mathcal{I}$ needs not to be an attractor. If $\Omega$ is regular, $f(x, u)$ is dissipative and $\mathcal{I}$ is the global attractor, we give an explicit bound on the Hausdorff dimension of $\mathcal{I}$ in terms of the structure parameter of the equation.

## 1. Introduction

In this paper we consider the reaction diffusion equation

$$
\begin{align*}
u_{t}+\beta(x) u-\Delta u & =f(x, u), & & (t, x) \in[0,+\infty[\times \Omega,  \tag{1.1}\\
u & =0, & & (t, x) \in[0,+\infty[\times \partial \Omega .
\end{align*}
$$

[^0]Here $\Omega$ is an arbitrary (possibly unbounded) open set in $\mathbb{R}^{3}, \beta(x)$ is a potential such that the operator $-\Delta+\beta(x)$ is positive, and $f(x, u)$ is a nonlinearity of subcritical growth (i.e. of polynomial growth strictly less than five).

The assumptions on $\beta(x)$ and $f(x, u)$ will be made more precise in Section 2 below. Under such assumptions, equation (1.1) generates a local semiflow $\pi$ in the space $H_{0}^{1}(\Omega)$. Suppose that the semiflow $\pi$ admits a compact invariant set $\mathcal{I}$ (i.e. $\pi(t, \mathcal{I})=\mathcal{I}$ for all $t \geq 0$ ). We do not make any structure assumption on the nonlinearity $f(x, u)$ and therefore we do not assume that $\mathcal{I}$ is the global attractor of equation (1.1): for example, $\mathcal{I}$ can be an unstable invariant set detected by Conley index arguments (see e.g. [16]).

Our aim is to prove that $\mathcal{I}$ has finite Hausdorff dimension and to give an explicit estimate of its dimension. The first results concerning the dimension of invariant sets of dynamical systems are due to Mallet-Paret [14] and Mañé [15]. For a comprehensive study of the subject, see e.g. [6], [12], [20], [23].

When $\Omega$ is a bounded domain and $f(x, u)$ satisfies suitable dissipativeness conditions, the existence of a finite dimensional compact global attractor for (1.1) is a classical achievement (see e.g. [6], [12], [23]). When $\Omega$ is unbounded, new difficulties arise due to the lack of compactness of the Sobolev embeddings. These difficulties can be overcome in several ways: by introducing weighted spaces (see e.g. [5], [9]), by developing suitable tail-estimates (see e.g. [24], [17]), by exploiting comparison arguments (see e.g. [3]).

Concerning the finite dimensionality of the attractor, in [5], [9], [24] and other similar works the potential $\beta(x)$ is always assumed to be just a positive constant. In [4] Arrieta et al. considered for the first time the case of a signchanging potential. In their results the invariant set $\mathcal{I}$ does not need to be an attractor; however they need to make some structure assumptions on $f(x, u)$ which essentially resemble the conditions ensuring the existence of the global attractor. Moreover, in [4] the invariant set is a-priori assumed to be bounded in the $L^{\infty}$-norm. In concrete situations, such a-priori estimate can be obtained through elliptic regularity combined with some comparison argument. This in turn requires to make some regularity assumption on the boundary of $\Omega$.

In this paper we do not make any structure assumption on the nonlinearity $f(x, u)$, neither do we assume $\partial \Omega$ to be regular. Our only assumption is that the mapping $h \mapsto\left(\partial_{u} f(x, 0)\right)_{+} h$ has to be a relatively form compact perturbation of $-\Delta+\beta(x)$. This can be achieved, e.g. by assuming that $\partial_{u} f(x, 0)$ can be estimated from above by some positive $L^{r}$ function, $r>3 / 2$. Under this assumption, we shall prove that $\mathcal{I}$ has finite Hausdorff dimension. Also, we give an explicit estimate of the dimension of $\mathcal{I}$, involving the number $\mathcal{N}$ of negative eigenvalues of the operator $-\Delta+\beta(x)-\partial_{u} f(x, 0)$. When $\Omega$ has a regular boundary, we can explicitly estimate $\mathcal{N}$ by mean of Cwickel-Lieb-Rozenblum inequality (see [21]);
as a consequence, if we also assume that $f(x, u)$ is dissipative, we recover the result of Arrieta et al. [4].

The paper is organized as follows. In Section 2 we introduce notations, we state the main assumptions and we collect some preliminaries about the semiflow generated by equation (1.1). In Section 3 we prove that the semiflow generated by equation (1.1) is uniformly $L^{2}$-differentiable on any compact invariant set $\mathcal{I}$. In Section 4 we recall the definition of Hausdorff dimension and we prove that any compact invariant set $\mathcal{I}$ has finite Hausdorff dimension in $L^{2}(\Omega)$ as well as in $H_{0}^{1}(\Omega)$. In Section 5 we compute the number of negative eigenvalues of the operator $-\Delta+\beta(x)-\partial_{u} f(x, 0)$ by mean of Cwickel-Lieb-Rozenblum inequality. In Section 6 we specialize our result to the case of a dissipative equation and we recover the result of Arrieta et al. [4].

The results contained in this paper continue to hold if one replaces $-\Delta$ with the general second order elliptic operator in divergence form

$$
-\sum_{i, j=1}^{3} \partial_{x_{i}}\left(a_{i j}(x) \partial_{x_{j}}\right)
$$

## 2. Notation, preliminaries and remarks

Let $\sigma \geq 1$. We denote by $L_{\mathrm{u}}^{\sigma}\left(\mathbb{R}^{N}\right)$ the set of measurable functions $\omega: \mathbb{R}^{N} \rightarrow \mathbb{R}$ such that

$$
|\omega|_{L_{\mathrm{u}}^{\sigma}}:=\sup _{y \in \mathbb{R}^{N}}\left(\int_{B(y)}|\omega(x)|^{\sigma} d x\right)^{1 / \sigma}<\infty
$$

where, for $y \in \mathbb{R}^{N}, B(y)$ is the open unit cube in $\mathbb{R}^{N}$ centered at $y$.
In this paper we assume throughout that $N=3$, and we fix an open (possibly unbounded) set $\Omega \subset \mathbb{R}^{3}$. We denote by $M_{B}$ the constant of the Sobolev embedding $H^{1}(B) \subset L^{6}(B)$, where $B$ is any open unit cube in $\mathbb{R}^{3}$. Moreover, for $2 \leq q \leq 6$, we denote by $M_{q}$ the constant of the Sobolev embedding $H^{1}\left(\mathbb{R}^{3}\right) \subset L^{q}\left(\mathbb{R}^{3}\right)$.

Proposition 2.1. Let $\sigma>3 / 2$ and let $\omega \in L_{\mathrm{u}}^{\sigma}\left(\mathbb{R}^{3}\right)$. Set $\rho:=3 / 2 \sigma$. Then, for every $\varepsilon>0$ and for every $u \in H_{0}^{1}(\Omega)$,

$$
\int_{\Omega}|\omega(x)||u(x)|^{2} d x \leq|\omega|_{L_{\mathrm{u}}^{\sigma}}\left(\rho \varepsilon M_{B}^{2}|u|_{H^{1}}^{2}+(1-\rho) \varepsilon^{-\rho /(1-\rho)}|u|_{L^{2}}^{2}\right) .
$$

Moreover, for every $u \in H_{0}^{1}(\Omega)$,

$$
\int_{\Omega}|\omega(x)||u(x)|^{2} d x \leq M_{B}^{2 \rho}|\omega|_{L_{\mathrm{u}}^{\sigma}}|u|_{H^{1}}^{2 \rho}|u|_{L^{2}}^{2(1-\rho)} .
$$

Proof. See the proof of Lemma 3.3 in [18].

Let $\beta \in L_{\mathrm{u}}^{\sigma}\left(\mathbb{R}^{3}\right)$, with $\sigma>3 / 2$. Let us consider the following bilinear form defined on the space $H_{0}^{1}(\Omega)$ :

$$
a(u, v):=\int_{\Omega} \nabla u(x) \cdot \nabla v(x) d x+\int_{\Omega} \beta(x) u(x) v(x) d x, \quad u, v \in H_{0}^{1}(\Omega)
$$

Our first assumption is the following:
Hypothesis 2.2. There exists $\lambda_{1}>0$ such that

$$
\begin{equation*}
\int_{\Omega}|\nabla u(x)|^{2} d x+\int_{\Omega} \beta(x)|u(x)|^{2} d x \geq \lambda_{1}|u|_{L^{2}}^{2}, \quad u \in H_{0}^{1}(\Omega) \tag{2.1}
\end{equation*}
$$

Remark 2.3. Conditions on $\beta(x)$ under which Hypothesis 2.2 is satisfied are expounded e.g. in [1], [2].

As a consequence of (2.1) and Proposition 2.1, we have:
Proposition 2.4. There exist two positive constants $\lambda_{0}$ and $\Lambda_{0}$ such that

$$
\lambda_{0}|u|_{H^{1}}^{2} \leq \int_{\Omega}|\nabla u(x)|^{2} d x+\int_{\Omega} \beta(x)|u(x)|^{2} d x \leq \Lambda_{0}|u|_{H^{1}}^{2}, \quad u \in H_{0}^{1}(\Omega) .
$$

The constants $\lambda_{0}$ and $\Lambda_{0}$ can be computed explicitly in terms of $\lambda_{1}, M_{B}$ and $|\beta|_{L_{u}^{\sigma}}$.
Proof. Cf. Lemma 4.2 in [17].
It follows from Proposition 2.4 that the bilinear form $a(\cdot, \cdot)$ defines a scalar product in $H_{0}^{1}(\Omega)$, equivalent to the standard one. According to the results of Section 4 in $[17], a(\cdot, \cdot)$ induces a positive selfadjoint operator $A$ in the space $L^{2}(\Omega)$. $A$ is uniquely determined by the relation

$$
\langle A u, v\rangle_{L^{2}}=a(u, v), \quad u \in D(A), v \in H_{0}^{1}(\Omega)
$$

Notice that $A u=-\Delta u+\beta u$ in the sense of distributions, and $u \in D(A)$ if and only if $-\Delta u+\beta u \in L^{2}(\Omega)$. Set $X:=L^{2}(\Omega)$, and let $\left(X^{\alpha}\right)_{\alpha \in \mathbb{R}}$ be the scale of fractional power spaces associated with $A$ (see Section 2 in [17] for a short, self-contained, description of this scale of spaces). Here we just recall that $X^{0}=L^{2}(\Omega), X^{1}=D(A), X^{1 / 2}=H_{0}^{1}(\Omega)$ and $X^{-\alpha}$ is the dual of $X^{\alpha}$ for $\alpha \in] 0,+\infty[$. For $\alpha \in] 0,+\infty\left[\right.$, the space $X^{\alpha}$ is a Hilbert space with respect to the scalar product

$$
\langle u, v\rangle_{X^{\alpha}}:=\left\langle A^{\alpha} u, A^{\alpha} v\right\rangle_{L^{2}}, \quad u, v \in X^{\alpha}
$$

Also, the space $X^{-\alpha}$ is a Hilbert space with respect to the scalar product $\langle\cdot, \cdot\rangle_{X^{-\alpha}}$ dual to the scalar product $\langle\cdot, \cdot\rangle_{X^{\alpha}}$, i.e.

$$
\left\langle u^{\prime}, v^{\prime}\right\rangle_{X^{-\alpha}}=\left\langle R_{\alpha}^{-1} u^{\prime}, R_{\alpha}^{-1} v^{\prime}\right\rangle_{X^{\alpha}}, \quad u, v \in X^{-\alpha}
$$

where $R_{\alpha}: X^{\alpha} \rightarrow X^{-\alpha}$ is the Riesz isomorphism $u \mapsto\langle\cdot, u\rangle_{X^{\alpha}}$. Finally, for every $\alpha \in \mathbb{R}, A$ induces a selfadjoint operator $A_{(\alpha)}: X^{\alpha+1} \rightarrow X^{\alpha}$, such that $A_{\left(\alpha^{\prime}\right)}$ is
an extension of $A_{(\alpha)}$ whenever $\alpha^{\prime} \leq \alpha$, and $D\left(A_{(\alpha)}^{\beta}\right)=X^{\alpha+\beta}$ for $\beta \in[0,1]$. If $\alpha \in[0,1 / 2], u \in X^{1-\alpha}$ and $v \in X^{1 / 2} \subset X^{\alpha}$, then

$$
\left\langle v, A_{(-\alpha)} u\right\rangle_{\left(X^{\alpha}, X^{-\alpha}\right)}=\langle u, v\rangle_{X^{1 / 2}}=a(u, v) .
$$

Lemma 2.5. Let $\left(X^{\alpha}\right)_{\alpha \in \mathbb{R}}$ be as above.
(a) If $p \in\left[2,6\left[\right.\right.$, then $X^{\alpha} \subset L^{p}(\Omega)$ for $\left.\left.\alpha \in\right] 3(p-2) / 4 p, 1 / 2\right]$. Accordingly, if $q \in] 6 / 5,2]$, then $L^{q}(\Omega) \subset X^{-\alpha}$ for $\left.\left.\alpha \in\right] 3(2-q) / 4 q, 1 / 2\right]$.
(b) If $\sigma>3 / 2$ and $\omega \in L_{\mathrm{u}}^{\sigma}(\Omega)$, then the assignment $u \mapsto \omega u$ defines a bounded linear map from $X^{1 / 2}$ to $X^{-\alpha}$ for $\left.\left.\alpha \in\right] 3 / 4 \sigma, 1 / 2\right]$.

Proof. See Lemmas 5.1 and 5.2 and the proof of Proposition 5.3 in [17].
Our second assumption is the following:
Hypothesis 2.6.
(a) $f: \Omega \times \mathbb{R} \rightarrow \mathbb{R}$ is such that, for every $u \in \mathbb{R}, f(\cdot, u)$ is measurable and, for almost every $x \in \Omega, f(x, \cdot)$ is of class $C^{2}$;
(b) $f(\cdot, 0) \in L^{q}(\Omega)$, with $6 / 5<q \leq 2$ and $\partial_{u} f(\cdot, 0) \in L_{\mathrm{u}}^{\sigma}\left(\mathbb{R}^{3}\right)$, with $\sigma>3 / 2$;
(c) there exist constants $C$ and $\gamma$, with $C>0$ and $2 \leq \gamma<3$ such that $\left|\partial_{u u} f(x, u)\right| \leq C\left(1+|u|^{\gamma}\right)$. Notice that, in view of Young's inequality, the requirement $\gamma \geq 2$ is not restrictive.

We introduce the Nemitski operator $\hat{f}$ which associates with every function $u: \Omega \rightarrow \mathbb{R}$ the function $\widehat{f}(u)(x):=f(x, u(x))$.

Proposition 2.7. Assume $f$ satisfies Hypothesis (2.6). Let $\alpha$ be such that

$$
\frac{1}{2}>\alpha>\max \left\{\frac{\gamma-1}{4}, \frac{3}{4} \frac{2-q}{q}, \frac{3}{4 \sigma}\right\}
$$

Then the assignment $u \mapsto \mathbf{f}(u)$, where

$$
\langle v, \mathbf{f}(u)\rangle_{\left(X^{\alpha}, X^{-\alpha}\right)}:=\int_{\Omega} \hat{f}(u)(x) v(x) d x
$$

defines a map $\mathbf{f}: X^{1 / 2} \rightarrow X^{-\alpha}$ which is Lipschitzian on bounded sets.
Proof. See the proof of Proposition 5.3 in [17].
Setting $\mathbf{X}:=X^{-\alpha}$ and $\mathbf{A}:=A_{(-\alpha)}$, we have that $\mathbf{X}^{\alpha+1 / 2}=X^{1 / 2}$. We can rewrite equation (1.1) as an abstract parabolic problem in the space $\mathbf{X}$, namely

$$
\begin{equation*}
\dot{u}+\mathbf{A} u=\mathbf{f}(u) . \tag{2.2}
\end{equation*}
$$

By results in [11], equation (2.2) has a unique mild solution for every initial datum $u_{0} \in \mathbf{X}^{\alpha+1 / 2}=H_{0}^{1}(\Omega)$, satisfying the variation of constants formula

$$
u(t)=e^{-\mathbf{A} t} u_{0}+\int_{0}^{t} e^{-\mathbf{A}(t-s)} \mathbf{f}(u(s)) d s, \quad t \geq 0
$$

It follows that (2.2) generates a local semiflow $\pi$ in the space $H_{0}^{1}(\Omega)$. Moreover, if $u(\cdot):\left[0, T\left[\rightarrow \mathbf{X}^{\alpha+1 / 2}\right.\right.$ is a mild solution of $(2.2)$, then $u(t)$ is differentiable into $\mathbf{X}^{\alpha+1 / 2}=H_{0}^{1}(\Omega)$ for $\left.t \in\right] 0, T[$, and it satisfies equation (2.2) in $\mathbf{X}=X^{-\alpha} \subset H^{-1}(\Omega)$. In particular, $u(\cdot)$ is a weak solution of (1.1).

Assume now that $\mathcal{I} \subset H_{0}^{1}(\Omega)$ is a compact invariant set for the semiflow $\pi$ generated by (2.2). If $\mathcal{B}$ is a Banach space such that $H_{0}^{1}(\Omega) \subset \mathcal{B}$, we define

$$
|\mathcal{I}|_{\mathcal{B}}:=\max \left\{|u|_{\mathcal{B}} \mid u \in \mathcal{I}\right\} .
$$

We end this section with a technical lemma that will be used later.
Lemma 2.8. For every $T>0$ there exists a constant $L(T)$ such that, whenever $u_{0}$ and $v_{0} \in \mathcal{I}$, setting $u(t):=\pi\left(t, u_{0}\right)$ and $v(t):=\pi\left(t, v_{0}\right), t \geq 0$, the following estimate holds:

$$
\left.\left.|u(t)-v(t)|_{H^{1}} \leq L(T) t^{-(\alpha+1 / 2)}\left|u_{0}-v_{0}\right|_{L^{2}}, \quad t \in\right] 0, T\right] .
$$

The constant $L(T)$ depends only on $|\mathcal{I}|_{H^{1}}$, and on the constants of Hypotheses 2.2 and 2.6.

Proof. We have

$$
u(t)-v(t)=e^{-\mathbf{A} t}\left(u_{0}-v_{0}\right)+\int_{0}^{t} e^{-\mathbf{A}(t-s)}(\mathbf{f}(u(s))-\mathbf{f}(v(s))) d s
$$

it follows that

$$
\begin{aligned}
& |u(t)-v(t)|_{\mathbf{X}^{\alpha+1 / 2}} \\
& \quad \leq t^{-(\alpha+1 / 2)}\left|u_{0}-v_{0}\right| \mathbf{x}+\int_{0}^{t}(t-s)^{-(\alpha+1 / 2)}|\mathbf{f}(u(s))-\mathbf{f}(v(s))| \mathbf{x} d s \\
& \quad \leq t^{-(\alpha+1 / 2)}\left|u_{0}-v_{0}\right| \mathbf{X}+\int_{0}^{t}(t-s)^{-(\alpha+1 / 2)} C\left(|\mathcal{I}|_{H^{1}}\right)|u(s)-v(s)|_{\mathbf{X}^{\alpha+1 / 2}} d s
\end{aligned}
$$

By Henry's inequality [11, Theorem 7.1.1], this implies that

$$
\left.\left.|u(t)-v(t)|_{\mathbf{x}^{\alpha+1 / 2}} \leq L(T) t^{-(\alpha+1 / 2)}\left|u_{0}-v_{0}\right| \mathbf{x}, \quad t \in\right] 0, T\right]
$$

and the thesis follows.

## 3. Uniform differentiability

In this section we prove some technical results which will allow us to apply the methods of [23] for proving finite dimensionality of compact invariant sets. We assume throughout that $\mathcal{I} \subset H_{0}^{1}(\Omega)$ is a compact invariant set of the semiflow $\pi$ generated by equation (2.2).

Lemma 3.1. There exists a constant $K$ such that, whenever $u_{0}$ and $v_{0} \in \mathcal{I}$, setting $u(t):=\pi\left(t, u_{0}\right)$ and $v(t):=\pi\left(t, v_{0}\right), t \geq 0$, the following estimate holds:

$$
|u(t)-v(t)|_{L^{2}}^{2}+\lambda_{0} \int_{0}^{t}|u(s)-v(s)|_{H^{1}}^{2} d s \leq e^{K t}\left|u_{0}-v_{0}\right|_{L^{2}}^{2}
$$

The constant $K$ depends only on $|\mathcal{I}|_{H^{1}}$, on $\lambda_{0}$ and $\Lambda_{0}$ (see Proposition 2.4), on $\left|\partial_{u} f(\cdot, 0)\right|_{L_{\mathrm{u}}^{\sigma}}$, and on the constants $C$ and $\gamma($ see Hypothesis 2.6).

Proof. Set $z(t)=u(t)-v(t)$. Then

$$
\begin{aligned}
\frac{1}{2} \frac{d}{d t}|z(t)|_{L^{2}}^{2}+\int_{\Omega}|\nabla z(t)(x)|^{2} d x & +\int_{\Omega} \beta(x)|z(t)(x)|^{2} d x \\
& =\int_{\Omega}(f(x, u(t)(x))-f(x, v(t)(x)) z(t)(x) d x
\end{aligned}
$$

It follows from Proposition 2.4 and Hypothesis 2.6 that

$$
\begin{aligned}
\frac{1}{2} \frac{d}{d t}|z(t)|_{L^{2}}^{2} & +\lambda_{0}|z(t)|_{H^{1}}^{2} \leq \int_{\Omega}\left|\partial_{u} f(x, 0)\right||z(t)(x)|^{2} d x \\
& +C^{\prime} \int_{\Omega}\left(1+|u(t)(x)|^{\gamma+1}+|v(t)(x)|^{\gamma+1}\right)|z(t)(x)|^{2} d x \\
\leq & \int_{\Omega}\left|\partial_{u} f(x, 0)\right||z(t)(x)|^{2} d x+C^{\prime}|z(t)|_{L^{2}}^{2} \\
& +C^{\prime}\left(|u(t)|_{L^{6}}^{\gamma+1}+|v(t)|_{L^{6}}^{\gamma+1}|z(t)|_{L^{12 /(5-\gamma)}}^{2},\right.
\end{aligned}
$$

where $C^{\prime}$ is a constant depending only on $C$ and $\gamma$. Notice that $2<12 /(5-\gamma)<6$. Therefore, by interpolation, we get that for every $\varepsilon>0$ there exists a constant $c_{\varepsilon}>0$ such that

$$
\begin{equation*}
|z(t)|_{L^{12 /(5-\gamma)}}^{2} \leq \varepsilon|z(t)|_{H^{1}}^{2}+c_{\varepsilon}|z(t)|_{L^{2}}^{2} . \tag{3.1}
\end{equation*}
$$

Now (3.1) and Proposition 2.1 imply that, for every $\varepsilon>0$, there exists a constant $C_{\varepsilon}^{\prime}$, depending on $C^{\prime},|\mathcal{I}|_{H^{1}}$ and $\varepsilon$, such that

$$
\begin{equation*}
\frac{1}{2} \frac{d}{d t}|z(t)|_{L^{2}}^{2}+\lambda_{0}|z(t)|_{H^{1}}^{2} \leq \varepsilon|z(t)|_{H^{1}}^{2}+C_{\varepsilon}^{\prime}|z(t)|_{L^{2}}^{2} \tag{3.2}
\end{equation*}
$$

Now choosing $\varepsilon=\lambda_{0} / 2$ and multiplying (3.2) by $e^{-2 C_{\varepsilon}^{\prime} t}$ we get

$$
\begin{equation*}
\frac{d}{d t}\left(e^{-2 C_{\varepsilon}^{\prime} t}|z(t)|_{L^{2}}^{2}\right)+\lambda_{0} e^{-2 C_{\varepsilon}^{\prime} t}|z(t)|_{H^{1}}^{2} \leq 0 \tag{3.3}
\end{equation*}
$$

Integrating (3.3) we obtain the thesis.

Let $\bar{u}(\cdot): \mathbb{R} \rightarrow H_{0}^{1}(\Omega)$ be a full bounded solution of $(2.2)$ such that $\bar{u}(t) \in \mathcal{I}$ for $t \in \mathbb{R}$. Let us consider the non autonomous linear equation

$$
\begin{array}{rlrl}
u_{t}+\beta(x) u-\Delta u & =\partial_{u} f(x, \bar{u}(t)) u,  \tag{3.4}\\
u & =0, & & (t, x) \in[0,+\infty[\times \Omega, \\
& & (t, x) \in[0,+\infty[\times \partial \Omega .
\end{array}
$$

We introduce the following bilinear form defined on on the space $H_{0}^{1}(\Omega)$ :

$$
\begin{align*}
a(t ; u, v):=\int_{\Omega} \nabla u(x) & \cdot \nabla v(x) d x+\int_{\Omega} \beta(x) u(x) v(x) d x  \tag{3.5}\\
& -\int_{\Omega} \partial_{u} f(x, \bar{u}(t)(x)) u(x) v(x) d x, \quad u, v \in H_{0}^{1}(\Omega)
\end{align*}
$$

Proposition 3.2. There exist constants $\kappa_{i}>0, i=1, \ldots, 4$, such that:
(a) $|a(t ; u, v)| \leq \kappa_{1}|u|_{H^{1}}|v|_{H^{1}}, u, v \in H_{0}^{1}(\Omega), t \in \mathbb{R}$;
(b) $|a(t ; u, u)| \geq \kappa_{2}|u|_{H^{1}}^{2}-\kappa_{3}|u|_{L^{2}}^{2}, u \in H_{0}^{1}(\Omega), t \in \mathbb{R}$;
(c) $|a(t ; u, v)-a(s ; u, v)| \leq \kappa_{4}|t-s||u|_{H^{1}}|v|_{H^{1}}, u, v \in H_{0}^{1}(\Omega), t, s \in \mathbb{R}$.

Proof. Properties (a) and (b) follow from Hypothesis 2.6 and Proposition 2.1. In order to prove point (c), we first observe that, by Theorem 3.5.2 in [11] (and its proof), the function $\bar{u}(\cdot)$ is differentiable into $H_{0}^{1}(\Omega)$, with $|\dot{\bar{u}}(\cdot)|_{H^{1}} \leq L$, where $L$ is a constant depending on $|\mathcal{I}|_{H^{1}}$ and on the constants in Hypotheses 2.2 and 2.6. Therefore we have:

$$
\begin{aligned}
\mid a(t ; u, v) & -a(s ; u, v)\left|\leq \int_{\Omega}\right| \partial_{u} f\left(x, \bar{u}(t)-\partial_{u} f(x, \bar{u}(s))| | u(x) \| v(x) \mid d x\right. \\
& \leq \int_{\Omega} C\left(1+|\bar{u}(t)(x)|^{\gamma}+|\bar{u}(s)(x)|^{\gamma}\right)|\bar{u}(t)(x)-\bar{u}(s)(x)||u(x) \| v(x)| d x \\
& \leq C^{\prime}\left(1+|\bar{u}(t)|_{H^{1}}^{\gamma}+|\bar{u}(s)|_{H^{1}}^{\gamma}\right)|\bar{u}(t)-\bar{u}(s)|_{H^{1}}|u|_{H^{1}}|v|_{H^{1}} \\
& \leq C^{\prime}\left(1+2|\mathcal{I}|_{H^{1}}^{\gamma}\right) L|t-s||u|_{H^{1}}|v|_{H^{1}},
\end{aligned}
$$

and the proof is complete.
Now let $A(t)$ be the self-adjoint operator determined by the relation

$$
\begin{equation*}
\langle A(t) u, v\rangle_{L^{2}}=a(t ; u, v), \quad u \in D(A(t)), v \in H_{0}^{1}(\Omega) . \tag{3.6}
\end{equation*}
$$

We can apply Theorem 3.1 in [10] and get:
Proposition 3.3. There exists a two parameter family of bounded linear operators $U(t, s): L^{2}(\Omega) \rightarrow L^{2}(\Omega), t \geq s$, such that:
(a) $U(s, s)=I$ for all $s \in \mathbb{R}$, and $U(t, s) U(s, r)=U(t, r)$ for all $t \geq s \geq r$;
(b) $U(t, s) h_{0} \in D(A(t))$ for all $h_{0} \in L^{2}(\Omega)$ and $t>s$;
(c) for every $h_{0} \in L^{2}(\Omega)$ and $s \in \mathbb{R}$, the map $t \mapsto U(t, s) h_{0}$ is differentiable into $L^{2}(\Omega)$ for $t>s$, and

$$
\frac{\partial}{\partial t} U(t, s) h_{0}=-A(t) U(t, s) h_{0}
$$

In particular, $U(t, s) h_{0}$ is a weak solution of (3.4).
Given $\bar{u}_{0} \in \mathcal{I}$, we take a full bounded solution $\bar{u}(\cdot)$ of (2.2), whose trajectory is contained in $\mathcal{I}$, and such that $\bar{u}(0)=\bar{u}_{0}$. Then we define

$$
\begin{equation*}
\mathcal{U}\left(\bar{u}_{0} ; t\right):=U(t, 0), \quad t \geq 0 \tag{3.7}
\end{equation*}
$$

where $U(t, s)$ is the family of operators given by Proposition 3.3. Notice that $\mathcal{U}\left(\bar{u}_{0} ; t\right)$ does not depend on the choice of $\bar{u}(\cdot)$, due to forward uniqueness for equation (2.2).

Proposition 3.4. For every $t \geq 0$,

$$
\sup _{\bar{u}_{0} \in \mathcal{I}}\left\|\mathcal{U}\left(\bar{u}_{0} ; t\right)\right\|_{\mathcal{L}\left(L^{2}, L^{2}\right)}<+\infty .
$$

Proof. Let $\bar{u}_{0} \in \mathcal{I}$ and $h_{0} \in L^{2}(\Omega)$. Set $h(t):=\mathcal{U}\left(\bar{u}_{0} ; t\right) h_{0}$. Then, by property (c) of Proposition 3.3, for $t>0$ we have

$$
\begin{aligned}
\frac{d}{d t} \frac{1}{2}|h(t)|_{L^{2}}^{2}+\int_{\Omega}|\nabla h(t)(x)|^{2} d x+\int_{\Omega} \beta(x) & |h(t)(x)|^{2} d x \\
& =\int_{\Omega} \partial_{u} f(x, \bar{u}(t)(x))|h(t)(x)|^{2} d x
\end{aligned}
$$

where $\bar{u}(\cdot)$ is a full bounded solution of (2.2), whose trajectory is contained in $\mathcal{I}$, and such that $\bar{u}(0)=\bar{u}_{0}$. It follows from Hypothesis 2.6 and Propositions 2.1 and 2.4 that for all $\varepsilon>0$

$$
\begin{aligned}
& \frac{d}{d t} \frac{1}{2}|h(t)|_{L^{2}}^{2}+\lambda_{0}|h(t)|_{H^{1}}^{2} \\
& \quad \leq \int_{\Omega} \partial_{u} f(x, 0)|h(t)(x)|^{2} d x+\int_{\Omega}\left(\partial_{u} f(x, \bar{u}(t)(x))-\partial_{u} f(x, 0)\right)|h(t)(x)|^{2} d x \\
& \quad \leq \varepsilon|h(t)|_{H^{1}}^{2}+c_{\varepsilon}|h(t)|_{L^{2}}^{2}+\int_{\Omega} C\left(1+|\bar{u}(t)(x)|^{\gamma}\right)|\bar{u}(t)(x)||h(t)(x)|^{2} d x \\
& \quad \leq \varepsilon|h(t)|_{H^{1}}^{2}+c_{\varepsilon}|h(t)|_{L^{2}}^{2}+\int_{\Omega} C^{\prime}\left(1+|\bar{u}(t)(x)|^{\gamma+1}\right)|h(t)(x)|^{2} d x \\
& \quad \leq \varepsilon|h(t)|_{H^{1}}^{2}+\left(c_{\varepsilon}+C^{\prime}\right)|h(t)|_{L^{2}}^{2}+C^{\prime}|\bar{u}(t)|_{L^{6}}^{\gamma+1}|h(t)|_{L^{12 /(5-\gamma)}}^{2} .
\end{aligned}
$$

Since $2<12 /(5-\gamma)<6$, by interpolation we get that for every $\varepsilon>0$ there exists a constant $c_{\varepsilon}^{\prime}>0$ such that

$$
|h(t)|_{L^{12 /(5-\gamma)}}^{2} \leq \varepsilon|h(t)|_{H^{1}}^{2}+c_{\varepsilon}^{\prime}|h(t)|_{L^{2}}^{2} .
$$

Therefore we have

$$
\begin{equation*}
\frac{d}{d t} \frac{1}{2}|h(t)|_{L^{2}}^{2}+\lambda_{0}|h(t)|_{H^{1}}^{2} \leq \varepsilon|h(t)|_{H^{1}}^{2}+C^{\prime \prime}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)|h(t)|_{L^{2}}^{2} \tag{3.8}
\end{equation*}
$$

Choosing $\varepsilon=\lambda_{0}$ and integrating (3.8) we obtain

$$
|h(t)|_{L^{2}}^{2} \leq e^{2 C^{\prime \prime}\left(\lambda_{0},|\mathcal{I}|_{H^{1}}\right) t}\left|h_{0}\right|_{L^{2}}^{2}
$$

Proposition 3.5. For every $t \geq 0$,

$$
\lim _{\varepsilon \rightarrow 0} \sup _{\substack{\bar{u}_{0}, \bar{v}_{0} \in \mathcal{I} \\ 0<\left|\bar{u}_{0}-\bar{v}_{0}\right|_{L^{2}}<\varepsilon}} \frac{\left|\pi\left(t, \bar{v}_{0}\right)-\pi\left(t, \bar{u}_{0}\right)-\mathcal{U}\left(\bar{u}_{0} ; t\right)\left(\bar{v}_{0}-\bar{u}_{0}\right)\right|_{L^{2}}}{\left|\bar{v}_{0}-\bar{u}_{0}\right|_{L^{2}}}=0
$$

Proof. Let $\bar{u}_{0}, \bar{v}_{0} \in \mathcal{I}$. Set $\bar{u}(t):=\pi\left(t, \bar{u}_{0}\right), \bar{v}(t):=\pi\left(t, \bar{v}_{0}\right)$ and $\theta(t):=$ $\bar{v}(t)-\bar{u}(t)-\mathcal{U}\left(\bar{u}_{0} ; t\right)\left(\bar{v}_{0}-\bar{u}_{0}\right), t \geq 0$. A computation using property (c) of Proposition 3.3 shows that, for $t>0$,

$$
\begin{aligned}
& \frac{d}{d t} \frac{1}{2}|\theta(t)|_{L^{2}}^{2}+\int_{\Omega}|\nabla \theta(t)(x)|^{2} d x+\int_{\Omega} \beta(x)|\theta(t)(x)|^{2} d x \\
& =\int_{\Omega} \partial_{u} f(x, \bar{u}(t)(x))|\theta(t)(x)|^{2} d x \\
& +\int_{\Omega}\left(f(x, \bar{v}(t)(x))-f(x, \bar{u}(t)(x))-\partial_{u} f(x, \bar{u}(t)(x))(\bar{v}(t)(x)-\bar{u}(t)(x))\right) \theta(t)(x) d x
\end{aligned}
$$

Therefore, by Proposition 2.4

$$
\frac{d}{d t} \frac{1}{2}|\theta(t)|_{L^{2}}^{2}+\lambda_{0}|\theta(t)|_{H^{1}} \leq I_{1}(t)+I_{2}(t)+I_{3}(t)
$$

where

$$
\begin{aligned}
I_{1}(t) & :=\int_{\Omega} \partial_{u} f(x, 0)|\theta(t)(x)|^{2} d x \\
I_{2}(t) & :=\int_{\Omega}\left(\partial_{u} f(x, \bar{u}(t)(x))-\partial_{u} f(x, 0)\right)|\theta(t)(x)|^{2} d x \\
I_{3}(t) & =\int_{\Omega}\left(f(x, \bar{v}(t))-f(x, \bar{u}(t))-\partial_{u} f(x, \bar{u}(t))(\bar{v}(t)-\bar{u}(t))\right) \theta(t) d x
\end{aligned}
$$

Repeating the same computations of the proof of Proposition 3.4, for $\varepsilon>0$ we get

$$
I_{1}(t)+I_{2}(t) \leq \varepsilon|\theta(t)|_{H^{1}}^{2}+C_{1}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)|\theta(t)|_{L^{2}}^{2}
$$

Concerning $I_{3}(t)$, for $\varepsilon>0$ we have

$$
\begin{aligned}
I_{2}(t) \leq & \int_{\Omega} C\left(1+|\bar{u}(t)(x)|^{\gamma}+|\bar{v}(t)(x)|^{\gamma}\right)|\bar{v}(t)(x)-\bar{u}(t)(x)|^{2} \theta(t)(x) d x \\
\leq & C|\theta(t)|_{L^{6}}|\bar{v}(t)-\bar{u}(t)|_{L^{12 / 5}}^{2}+C|\theta(t)|_{L^{6}}\left(|\bar{u}(t)|_{L^{6}}^{\gamma}\right. \\
& \left.+|\bar{v}(t)|_{L^{6}}^{\gamma}\right)|\bar{v}(t)-\bar{u}(t)|_{L^{12 /(5-\gamma)}}^{2} \\
\leq & \varepsilon|\theta(t)|_{H^{1}}^{2}+C_{2}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)\left(|\bar{v}(t)-\bar{u}(t)|_{L^{12 / 5}}^{4}+|\bar{v}(t)-\bar{u}(t)|_{L^{12 /(5-\gamma)}}^{4}\right) \\
\leq & \varepsilon|\theta(t)|_{H^{1}}^{2}+C_{3}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)\left(|\bar{v}(t)-\bar{u}(t)|_{H^{1}}|\bar{v}(t)-\bar{u}(t)|_{L^{2}}^{3}\right. \\
& \left.+|\bar{v}(t)-\bar{u}(t)|_{H^{1}}^{1+\gamma}|\bar{v}(t)-\bar{u}(t)|_{L^{2}}^{3-\gamma}\right)
\end{aligned}
$$

Choosing $\varepsilon=\lambda_{0} / 2$, we get

$$
\begin{aligned}
& \quad \frac{d}{d t} \frac{1}{2}|\theta(t)|_{L^{2}}^{2}-C_{1}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)|\theta(t)|_{L^{2}}^{2} \\
& \leq C_{3}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)\left(|\bar{v}(t)-\bar{u}(t)|_{H^{1}}|\bar{v}(t)-\bar{u}(t)|_{L^{2}}^{3}+|\bar{v}(t)-\bar{u}(t)|_{H^{1}}^{1+\gamma}|\bar{v}(t)-\bar{u}(t)|_{L^{2}}^{3-\gamma}\right) \\
& \\
& \quad \leq C_{4}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)\left(|\bar{v}(t)-\bar{u}(t)|_{L^{2}}^{3}+|\bar{v}(t)-\bar{u}(t)|_{H^{1}}^{2}|\bar{v}(t)-\bar{u}(t)|_{L^{2}}^{3-\gamma}\right)
\end{aligned}
$$

By Lemma 3.1, we get

$$
\begin{aligned}
& \frac{d}{d t} \frac{1}{2}|\theta(t)|_{L^{2}}^{2}-C_{1}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)|\theta(t)|_{L^{2}}^{2} \\
& \quad \leq C_{4}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)\left(e^{3 K t}\left|\bar{v}_{0}-\bar{u}_{0}\right|_{L^{2}}^{3}+e^{(3-\gamma) K t}|\bar{v}(t)-\bar{u}(t)|_{H^{1}}^{2}\left|\bar{v}_{0}-\bar{u}_{0}\right|_{L^{2}}^{3-\gamma}\right)
\end{aligned}
$$

Writing $C_{1}$ for $C_{1}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)$ and $C_{4}$ for $C_{4}\left(\varepsilon,|\mathcal{I}|_{H^{1}}\right)$, we have

$$
\begin{align*}
& \frac{d}{d t} \frac{1}{2}\left(e^{-C_{1} t}|\theta(t)|_{L^{2}}^{2}\right)  \tag{3.9}\\
& \leq C_{4}\left(e^{\left(3 K-C_{1}\right) t}\left|\bar{v}_{0}-\bar{u}_{0}\right|_{L^{2}}^{3}+e^{\left((3-\gamma) K-C_{1}\right) t}|\bar{v}(t)-\bar{u}(t)|_{H^{1}}^{2}\left|\bar{v}_{0}-\bar{u}_{0}\right|_{L^{2}}^{3-\gamma}\right)
\end{align*}
$$

Finally, integrating (3.9), recalling that $\theta(0)=0$ and taking into account Lemma 3.1, we get the existence of two increasing functions $\Phi_{1}(t)$ and $\Phi_{2}(t)$ such that

$$
|\theta(t)|_{L^{2}}^{2} \leq \Phi_{1}(t)\left|\bar{v}_{0}-\bar{u}_{0}\right|_{L^{2}}^{3}+\Phi_{2}(t)\left|\bar{v}_{0}-\bar{u}_{0}\right|_{L^{2}}^{5-\gamma}
$$

and the thesis follows.

## 4. Dimension of invariant sets

Let $\mathcal{X}$ be a complete metric space and let $\mathcal{K} \subset \mathcal{X}$ be a compact set. For $d \in \mathbb{R}^{+}$and $\varepsilon>0$ one defines

$$
\mu_{H}(\mathcal{K}, d, \varepsilon):=\inf \left\{\sum_{i \in I} r_{i}^{d} \mid \mathcal{K} \subset \bigcup_{i \in I} B\left(x_{i}, r_{i}\right), r_{i} \leq \varepsilon\right\}
$$

where the infimum is taken over all the finite coverings of $\mathcal{K}$ with balls of radius $r_{i} \leq \varepsilon$. Observe that $\mu_{H}(\mathcal{K}, d, \varepsilon)$ is a non increasing function of $\varepsilon$ and $d$. The $d$-dimensional Hausdorff measure of $\mathcal{K}$ is by definition

$$
\mu_{H}(\mathcal{K}, d):=\lim _{\varepsilon \rightarrow 0} \mu_{H}(\mathcal{K}, d, \varepsilon)=\sup _{\varepsilon>0} \mu_{H}(\mathcal{K}, d, \varepsilon)
$$

One has:
(1) $\mu_{H}(\mathcal{K}, d) \in[0,+\infty]$;
(2) if $\mu_{H}(\mathcal{K}, \bar{d})<\infty$, then $\mu_{H}(\mathcal{K}, d)=0$ for all $d>\bar{d}$;
(3) if $\mu_{H}(\mathcal{K}, \bar{d})>0$, then $\mu_{H}(\mathcal{K}, d)=+\infty$ for all $d<\bar{d}$.

The Hausdorff dimension of $\mathcal{K}$ is the smallest $d$ for which $\mu_{H}(\mathcal{K}, d)$ is finite, i.e.

$$
\operatorname{dim}_{H}(\mathcal{K}):=\inf \left\{d>0 \mid \mu_{H}(\mathcal{K}, d)=0\right\} .
$$

As pointed up in [22], the Hausdorff dimension is in fact an intrinsic metric property of the set $\mathcal{K}$. Moreover, if $\mathcal{Y}$ is another complete metric space and $\ell: \mathcal{K} \rightarrow \mathcal{Y}$ is a Lipschitzian map, then $\operatorname{dim}_{H}(\ell(\mathcal{K})) \leq \operatorname{dim}_{H}(\mathcal{K})$.

There is a well developed technique to estimate the Hausdorff dimension of an invariant set of a map or a semigroup. We refer the reader e.g. to [23] and [12]. The geometric idea consists in tracking the evolution of a $d$-dimensional volume under the action of the linearization of the semigroup along solutions lying in the invariant set. One looks then for the smallest $d$ for which any $d$-dimensional volume contracts asymptotically as $t \rightarrow \infty$.

Let $\bar{u}_{0} \in \mathcal{I}$ and let $\bar{u}(\cdot): \mathbb{R} \rightarrow H_{0}^{1}(\Omega)$ be a full bounded solution of (2.2) such that $\bar{u}(0)=\bar{u}_{0}$ and $\bar{u}(t) \in \mathcal{I}$ for $t \in \mathbb{R}$. For $t \geq 0$, we denote by $a_{\bar{u}_{0}}(t ; u, v)$ the bilinear form defined by (3.5), and by $A_{\bar{u}_{0}}(t)$ the self-adjoint operator determined by the relation (3.6). Given a $d$-dimensional subspace $E_{d}$ of $L^{2}(\Omega)$, with $E_{d} \subset$ $H_{0}^{1}(\Omega)$, we define the operator $A_{\bar{u}_{0}}\left(t \mid E_{d}\right): E_{d} \rightarrow E_{d}$ by

$$
\left\langle A_{\bar{u}_{0}}\left(t \mid E_{d}\right) \phi, \psi\right\rangle_{L^{2}}:=a_{\bar{u}_{0}}(t ; \phi, \psi), \quad \phi, \psi \in E_{d} .
$$

Notice that, if $E_{d} \subset D\left(A_{\bar{u}_{0}}(t)\right)$, then one has $A_{\bar{u}_{0}}\left(t \mid E_{d}\right)=P_{E_{d}} A_{\bar{u}_{0}}(t) P_{E_{d}} \mid E_{d}$, where $P_{E_{d}}: L^{2}(\Omega) \rightarrow E_{d}$ is the $L^{2}$-orthogonal projection onto $E_{d}$. We define

$$
\operatorname{Tr}_{d}\left(A_{\bar{u}_{0}}(t)\right):=\inf _{\substack{E_{d} \subset H_{0}^{1}(\Omega) \\ \operatorname{dim} E_{d}=d}} \operatorname{Tr}\left(A_{\bar{u}_{0}}\left(t \mid E_{d}\right)\right) .
$$

Let $\bar{u}_{0} \in \mathcal{I}$, let $d \in \mathbb{N}$ and let $v_{0, i} \in L^{2}(\Omega), i=1, \ldots, d$. Set $v_{i}(t):=$ $\mathcal{U}\left(\bar{u}_{0} ; t\right) v_{0, i}, t \geq 0$, where $\mathcal{U}\left(\bar{u}_{0} ; t\right)$ is defined by (3.7). We denote by $G(t)$ the $d$-dimensional volume delimited by $v_{1}(t), \ldots, v_{d}(t)$ in $L^{2}(\Omega)$, that is

$$
G(t):=\left|v_{1}(t) \wedge v_{2}(t) \wedge \cdots \wedge v_{d}(t)\right|_{\wedge^{d} L^{2}}=\left(\operatorname{det}\left(\left\langle v_{i}(t), v_{j}(t)\right\rangle_{L^{2}}\right)_{i j}\right)^{1 / 2}
$$

An easy computation using Leibnitz rule and Proposition 3.3 shows that, for $t>0, G(t)$ satisfies the ordinary differential equation

$$
G^{\prime}(t)=-\operatorname{Tr}\left(A_{\bar{u}_{0}}\left(t \mid E_{d}(t)\right) G(t)\right.
$$

where $E_{d}(t):=\operatorname{span}\left(v_{1}(t), \ldots, v_{d}(t)\right)$. It follows from Propositions 3.4 and 3.5 and from the results in [23, Chapter V] that the Hausdorff dimension $\operatorname{dim}_{H}(\mathcal{I})$ of $\mathcal{I}$ in $L^{2}(\Omega)$ is finite and less than or equal to $d$, provided

$$
\limsup _{t \rightarrow \infty} \sup _{\bar{u}_{0} \in \mathcal{I}} \frac{1}{t} \int_{0}^{t}-\operatorname{Tr}_{d}\left(A_{\bar{u}_{0}}(s)\right) d s<0 .
$$

In order to prove that $\operatorname{dim}_{H}(\mathcal{I}) \leq d$, we are lead to estimate $-\operatorname{Tr}_{d}\left(A_{\bar{u}_{0}}(t)\right)$. To this end, we notice that, whenever $E_{d}$ is a $d$-dimensional subspace of $L^{2}(\Omega)$, and $B: E_{d} \rightarrow E_{d}$ is a selfadjoint operator, then

$$
\operatorname{Tr}(B)=\sum_{i=1}^{d}\left\langle B \phi_{i}, \phi_{i}\right\rangle_{L^{2}},
$$

where $\phi_{1}, \ldots, \phi_{d}$ is any $L^{2}$-orthonormal basis of $E_{d}$. So let $E_{d} \subset H_{0}^{1}(\Omega)$ be a $d$-dimensional space and let $\phi_{1}, \ldots, \phi_{d}$ be an $L^{2}$-orthonormal basis of $E_{d}$. Fix $0<\delta<1$. It follows that

$$
\begin{aligned}
& \operatorname{Tr}\left(A_{\bar{u}_{0}}\left(t \mid E_{d}\right)\right) \\
&= \sum_{i=1}^{d}\left((1-\delta)\left(\int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x+\int_{\Omega} \beta(x)\left|\phi_{i}\right|^{2} d x\right)-\int_{\Omega} \partial_{u} f(x, 0)\left|\phi_{i}\right|^{2} d x\right) \\
& \quad+\delta \sum_{i=1}^{d}\left(\int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x+\int_{\Omega} \beta(x)\left|\phi_{i}\right|^{2} d x\right) \\
& \quad+\sum_{i=1}^{d} \int_{\Omega}\left(\partial_{u} f(x, \bar{u}(t))-\partial_{u} f(x, 0)\right)\left|\phi_{i}\right|^{2} d x .
\end{aligned}
$$

We introduce the following bilinear form defined on the space $H_{0}^{1}(\Omega)$ :

$$
\begin{aligned}
& a_{\delta}(u, v):=(1-\delta)\left(\int_{\Omega} \nabla u(x) \cdot \nabla v(x) d x+\int_{\Omega} \beta(x) u(x) v(x) d x\right) \\
&-\int_{\Omega} \partial_{u} f(x, 0) u(x) v(x) d x, \quad u, v \in H_{0}^{1}(\Omega)
\end{aligned}
$$

Let $A_{\delta}$ be the self-adjoint operator determined by the relation

$$
\left\langle A_{\delta} u, v\right\rangle_{L^{2}}=a_{\delta}(u, v), \quad u \in D\left(A_{\delta}\right), v \in H_{0}^{1}(\Omega)
$$

Given a $d$-dimensional subspace $E_{d}$ of $L^{2}(\Omega)$, with $E_{d} \subset H_{0}^{1}(\Omega)$, we define the operator $A_{\delta}\left(E_{d}\right): E_{d} \rightarrow E_{d}$ by

$$
\left\langle A_{\delta}\left(E_{d}\right) \phi, \psi\right\rangle_{L^{2}}:=a_{\delta}(\phi, \psi), \quad \phi, \psi \in E_{d}
$$

It follows that

$$
\begin{aligned}
\operatorname{Tr}\left(A_{\bar{u}_{0}}\left(t \mid E_{d}\right)\right) & =\operatorname{Tr}\left(A_{\delta}\left(E_{d}\right)\right)+\delta \sum_{i=1}^{d}\left(\int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x\right. \\
+ & \left.\int_{\Omega} \beta(x)\left|\phi_{i}\right|^{2} d x\right)+\sum_{i=1}^{d} \int_{\Omega}\left(\partial_{u} f(x, \bar{u}(t))-\partial_{u} f(x, 0)\right)\left|\phi_{i}\right|^{2} d x
\end{aligned}
$$

We introduce the proper values of the operator $A_{\delta}$ :

$$
\mu_{j}\left(A_{\delta}\right):=\sup _{\substack{\psi_{1}, \ldots, \psi_{j-1} \in H_{0}^{1}(\Omega)}} \inf _{\substack{\psi \in\left[\psi_{1}, \ldots, \psi_{j-1}\right]^{\perp} \\|\psi|_{L^{2}}=1, \psi \in H_{0}^{1}(\Omega)}} a_{\delta}(\psi, \psi) \quad j=1,2, \ldots
$$

We recall (see e.g. Theorem XIII. 1 in [19]) that:
Proposition 4.1. For each fixed $n$, either
(a) there are at least $n$ eigenvalues (counting multiplicity) below the bottom of the essential spectrum of $A_{\delta}$ and $\mu_{n}\left(A_{\delta}\right)$ is the $n$th eigenvalue (counting multiplicity);
(b) $\mu_{n}\left(A_{\delta}\right)$ is the bottom of the essential spectrum and in that case $\mu_{n+j}\left(A_{\delta}\right)$ $=\mu_{n}\left(A_{\delta}\right), j=1,2, \ldots$ and there are at most $n-1$ eigenvalues (counting multiplicity) below $\mu_{n}\left(A_{\delta}\right)$.

Let $\mu_{j}\left(A_{\delta}\left(E_{d}\right)\right), j=1, \ldots, d$, be the eigenvalues of $A_{\delta}\left(E_{d}\right)$. By Theorem XIII. 3 in [19], we have that

$$
\mu_{j}\left(A_{\delta}\left(E_{d}\right)\right) \geq \mu_{j}\left(A_{\delta}\right), \quad j=1, \ldots, d
$$

It follows that

$$
\begin{aligned}
\operatorname{Tr}\left(A_{\bar{u}_{0}}\left(t \mid E_{d}\right)\right) \geq \sum_{i=1}^{d} \mu_{i}\left(A_{\delta}\right)+ & \delta \sum_{i=1}^{d}\left(\int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x+\int_{\Omega} \beta(x)\left|\phi_{i}\right|^{2} d x\right) \\
& +\sum_{i=1}^{d} \int_{\Omega}\left(\partial_{u} f(x, \bar{u}(t))-\partial_{u} f(x, 0)\right)\left|\phi_{i}\right|^{2} d x .
\end{aligned}
$$

To proceed further, we need to recall the Lieb-Thirring inequality (see [13]).
Proposition 4.2. Let $N \in \mathbb{N}$ and let $p \in \mathbb{R}$, with $\max \{N / 2,1\} \leq p \leq$ $1+N / 2$. There exists a constant $K_{p, N}>0$ such that, if $\phi_{1}, \ldots, \phi_{d} \in H^{1}\left(\mathbb{R}^{N}\right)$ are pairwise $L^{2}$-orthonormal, then

$$
\begin{equation*}
\sum_{i=1}^{d} \int_{\mathbb{R}^{N}}\left|\nabla \phi_{i}(x)\right|^{2} d x \geq \frac{1}{K_{p, N}}\left(\int_{\mathbb{R}^{N}} \rho(x)^{p /(p-1)} d x\right)^{2(p-1) / N} \tag{4.1}
\end{equation*}
$$

where $\rho(x):=\sum_{i=1}^{d}\left|\phi_{i}(x)\right|^{2}$.
Now we have:
Lemma 4.3. Let $\bar{u} \in \mathcal{I}$ and let $\phi_{1}, \ldots, \phi_{d} \in H^{1}\left(\mathbb{R}^{N}\right)$ be pairwise $L^{2}$-orthonormal. Then

$$
\begin{aligned}
& \delta \sum_{i=1}^{d}\left(\int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x+\int_{\Omega} \beta(x)\left|\phi_{i}\right|^{2} d x\right) \\
& \quad+\sum_{i=1}^{d} \int_{\Omega}\left(\partial_{u} f(x, \bar{u}(x))-\partial_{u} f(x, 0)\right)\left|\phi_{i}\right|^{2} d x \geq-D\left(\gamma, \lambda_{0}, \delta,|\mathcal{I}|_{H^{1}}\right),
\end{aligned}
$$

where

$$
\begin{aligned}
D\left(\gamma, \lambda_{0}, \delta,|\mathcal{I}|_{H^{1}}\right) & =\frac{5}{2}\left(\frac{3}{5} \frac{2}{\delta \lambda_{0}}\right)^{3 / 2}\left(C|\mathcal{I}|_{L^{5 / 2}} K_{5 / 2,3}\right)^{5 / 2} \\
+ & \frac{3-\gamma}{4}\left(\frac{\gamma+1}{4} \frac{2}{\delta \lambda_{0}}\right)^{(\gamma+1) /(3-\gamma)}\left(C|\mathcal{I}|_{L^{6}}^{\gamma+1} K_{6 /(\gamma+1), 3}\right)^{4 /(3-\gamma)} .
\end{aligned}
$$

Proof. We observe first that

$$
\delta \sum_{i=1}^{d}\left(\int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x+\int_{\Omega} \beta(x)\left|\phi_{i}\right|^{2} d x\right) \geq \delta \lambda_{0} \sum_{i=1}^{d} \int_{\Omega}\left|\nabla \phi_{i}\right|^{2} d x .
$$

On the other hand,

$$
\begin{aligned}
& \left|\int_{\Omega}\left(\partial_{u} f(x, \bar{u}(x))-\partial_{u} f(x, 0)\right) \rho(x) d x\right| \\
& \quad \leq \int_{\Omega} C\left(1+|\bar{u}|^{\gamma}\right)|\bar{u}||\rho| d x \leq C|\bar{u}|_{L^{5 / 2}}|\rho|_{L^{5 / 3}}+C|\bar{u}|_{L^{6}}^{\gamma+1}|\rho|_{L^{6 /(5-\gamma)}} .
\end{aligned}
$$

By Lieb-Thirring inequality (4.1), we have

$$
\begin{aligned}
& \left|\int_{\Omega}\left(\partial_{u} f(x, \bar{u}(x))-\partial_{u} f(x, 0)\right) \rho(x) d x\right| \\
& \leq C|\mathcal{I}|_{L^{5 / 2}} K_{5 / 2,3}\left(\sum_{i=1}^{d} \int_{\mathbb{R}^{N}}\left|\nabla \phi_{i}\right|^{2} d x\right)^{3 / 5} \\
& \quad+C|\mathcal{I}|_{L^{6}}^{\gamma+1} K_{6 /(\gamma+1), 3}\left(\sum_{i=1}^{d} \int_{\mathbb{R}^{N}}\left|\nabla \phi_{i}\right|^{2} d x\right)^{(\gamma+1) / 4} .
\end{aligned}
$$

The conclusion follows by a simple application of Young's inequality.
Thanks to Lemma 4.3, we finally get:

$$
\operatorname{Tr}\left(A_{\bar{u}_{0}}\left(t \mid E_{d}\right)\right) \geq \sum_{i=1}^{d} \mu_{i}\left(A_{\delta}\right)-D\left(\gamma, \lambda_{0}, \delta,|\mathcal{I}|_{H^{1}}\right)
$$

Therefore, in order to conclude that $\operatorname{dim}_{H}(\mathcal{I})$ is finite, we are lead to make some assumption which guarantees that $\sum_{i=1}^{d} \mu_{i}\left(A_{\delta}\right)$ can be made positive and as large as we want, by choosing $d$ sufficiently large. This is equivalent to the fact that the bottom of the essential spectrum of $A_{\delta}$ be strictly positive. We make the following assumption:

Hypothesis 4.4. For every $\varepsilon>0$ there exists $V_{\varepsilon} \in L^{r}(\Omega), r>3 / 2, V_{\varepsilon} \geq 0$, such that $\partial_{u} f(x, 0) \leq V_{\varepsilon}(x)+\varepsilon$, for $x \in \Omega$.

We need the following lemmas:
Lemma 4.5. Let $r>3 / 2$ and let $V \in L^{r}(\Omega)$. If $r>3$ let $p:=2$; if $r \leq 3$ let $p:=6 / 5$. Then the assignment $u \mapsto V u$ defines a compact map from $H_{0}^{1}(\Omega)$ to $L^{p}(\Omega)$, and hence to $H^{-1}(\Omega)$.

Proof. Let $B \subset H_{0}^{1}(\Omega)$ be bounded. If $\mathcal{B}$ is a Banach space such that $H_{0}^{1}(\Omega) \subset \mathcal{B}$, we define $|B|_{\mathcal{B}}:=\sup \left\{|u|_{\mathcal{B}} \mid u \in B\right\}$. If $u \in H_{0}^{1}(\Omega)$ we denote by $\widetilde{u}$ its trivial extension to the whole $\mathbb{R}^{3}$. Similarly, we denote by $\widetilde{V}$ the trivial
extension of $V$ to $\mathbb{R}^{3}$. For $k>0$, let $\chi_{k}$ be the characteristic function of the set $\left\{x \in \mathbb{R}^{3}| | x \mid \leq k\right\}$. Now, for $u \in B$ and $k>0$, we have:

$$
\int_{\mathbb{R}^{3}}\left|\left(1-\chi_{k}\right) \widetilde{V} \widetilde{u}\right|^{p} d x \leq\left(\int_{|x| \geq k}|\widetilde{V}|^{r} d x\right)^{p / r}\left(\int_{|x| \geq k}|\widetilde{u}|^{p r /(r-p)} d x\right)^{(r-p) / r}
$$

It follows that

$$
\left|\left(1-\chi_{k}\right) \widetilde{V} \widetilde{u}\right|_{L^{p}} \leq|B|_{L^{p r /(r-p)}}\left|\left(1-\chi_{k}\right) \widetilde{V}\right|_{L^{r}}, \quad u \in B, k>0
$$

Similarly, we have:

$$
\begin{equation*}
\left|\chi_{k} \widetilde{V} \widetilde{u}\right|_{L^{p}} \leq|\widetilde{V}|_{L^{r}}\left|\chi_{k} \widetilde{u}\right|_{L^{p r /(r-p)}}, \quad u \in H_{0}^{1}(\Omega), k>0 . \tag{4.2}
\end{equation*}
$$

Now, given $\varepsilon>0$, we choose $k>0$ so large that $\left|\left(1-\chi_{k}\right) \widetilde{V}\right|_{L^{r}} \leq \varepsilon$. Then

$$
\begin{aligned}
\{\widetilde{V} \widetilde{u} \mid u \in B\} & =\left\{\chi_{k} \widetilde{V} \widetilde{u}+\left(1-\chi_{k}\right) \widetilde{V} \widetilde{u} \mid u \in B\right\} \\
& \subset\left\{\chi_{k} \widetilde{V} \widetilde{u} \mid u \in B\right\}+\left\{\left(1-\chi_{k}\right) \widetilde{V} \widetilde{u} \mid u \in B\right\} \\
& \subset\left\{\left.v \in L^{p}\left(\mathbb{R}^{3}\right)| | v\right|_{L^{p}} \leq \varepsilon\right\}+\left\{\chi_{k} \widetilde{V} \widetilde{u} \mid u \in B\right\} .
\end{aligned}
$$

We notice that $2 \leq p r /(r-p)<6$ : therefore, By Rellich's Theorem, $H^{1}\left(B_{k}(0)\right)$ is compactly embedded into $L^{p r /(r-p)}$. It follows that the set $\left\{\chi_{k} \widetilde{u} \mid u \in B\right\}$ is precompact in $L^{p r /(r-p)}$. By (4.2), we deduce that $\left\{\chi_{k} \widetilde{V} \widetilde{u} \mid u \in B\right\}$ is precompact in $L^{p}\left(\mathbb{R}^{3}\right)$. A simple measure of non compactness argument shows then that the set $\{\tilde{V} \widetilde{u} \mid u \in B\}$ is precompact in $L^{p}\left(\mathbb{R}^{3}\right)$ and this in turn implies that the set $\{V u \mid u \in B\}$ is precompact in $L^{p}(\Omega)$.

Lemma 4.6. Let $V$ be as in Lemma 4.5. Let $A+V$ be the selfadjoint operator determined by the bilinear form $a(u, v)+\int_{\Omega} V u v d x, u, v \in H_{0}^{1}(\Omega)$. Then, for sufficiently large $\lambda>0,(A+\lambda)^{-1}-(A+V+\lambda)^{-1}$ is a compact operator in $L^{2}(\Omega)$.

Proof. Take $\lambda>0$ so large that $A+V+\lambda$ be strictly positive. Let $u \in$ $L^{2}(\Omega)$. Set $v:=(A+V+\lambda)^{-1} u, w:=(A+\lambda)^{-1} u$ and $z:=v-w$. This means that

$$
a(v, \phi)+\lambda(v, \phi)+\int_{\Omega} V v \phi d x=\int_{\Omega} u \phi d x, \quad \text { for all } \phi \in H_{0}^{1}(\Omega)
$$

and

$$
a(w, \phi)+\lambda(w, \phi)=\int_{\Omega} u \phi d x, \quad \text { for all } \phi \in H_{0}^{1}(\Omega)
$$

It follows that

$$
a(z, \phi)+\lambda(z, \phi)+\int_{\Omega} V v \phi d x=0, \quad \text { for all } \phi \in H_{0}^{1}(\Omega)
$$

Choosing $\phi:=z$, Proposition 2.4 and Lemma 4.5 imply

$$
\lambda_{0}|z|_{H^{1}}^{2} \leq|z|_{H^{1}}|V v|_{H^{-1}} \leq \frac{\lambda_{0}}{2}|z|_{H^{1}}^{2}+K_{\lambda_{0}}|V v|_{H^{-1}}^{2}
$$

Therefore we obtain the estimate

$$
\left|(A+\lambda)^{-1} u-(A+V+\lambda)^{-1} u\right|_{H^{1}} \leq K_{\lambda_{0}}\left|V(A+V+\lambda)^{-1} u\right|_{H^{-1}}, \quad u \in L^{2}(\Omega)
$$ and the conclusion follows from Lemma 4.5.

Now we can prove:
Proposition 4.7. Assume Hypothesis 4.4 is satisfied. Then the essential spectrum of $A_{\delta}$ is contained in $\left[(1-\delta) \lambda_{1},+\infty[\right.$.

Proof. Hypothesis 4.4 and Proposition 4.1 imply that, for every $\varepsilon>0$, the bottom of the essential spectrum of $A_{\delta}$ is larger than or equal to the bottom of the essential spectrum of $(1-\delta) A-\varepsilon-V_{\varepsilon}(x)$. We observe that the spectrum of $(1-\delta) A-\varepsilon$ is contained in [(1- $\delta) \lambda_{1}-\varepsilon,+\infty[$. By Lemma 4.6 and Weyl's Theorem (see [19, Theorem XIII.14]), the essential spectrum of ( $1-\delta$ ) $A-\varepsilon-V_{\varepsilon}(x)$ coincides with that of $(1-\delta) A-\varepsilon$. It follows that the bottom of the essential spectrum of $A_{\delta}$ is larger than or equal to $(1-\delta) \lambda_{1}-\varepsilon$ for arbitrary small $\varepsilon>0$, and the conclusion follows.

Whenever Hypothesis 4.4 is satisfied, for $0<\delta<1$ and $\lambda<(1-\delta) \lambda_{0}$ we introduce the following quantity:

$$
\mathcal{N}(\delta, \lambda):=\# \text { eigenvalues of } A_{\delta} \text { below } \lambda
$$

Then, for $d \geq \mathcal{N}\left(\delta,\left(\frac{1-\delta}{2} \lambda_{1}\right)\right.$ we have:

$$
\sum_{i=1}^{d} \mu_{i}\left(A_{\delta}\right) \geq \mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right) \mu_{1}(\delta)+\left(d-\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right)\right) \frac{1-\delta}{2} \lambda_{1}
$$

We have thus proved our first main result:
Theorem 4.8. Assume Hypotheses 2.2, 2.6 and 4.4 are satisfied. Let $\mathcal{I} \subset$ $H_{0}^{1}(\Omega)$ be a compact invariant set for the semiflow $\pi$ generated by equation (2.2) in $H_{0}^{1}(\Omega)$. Then the Hausdorff dimension of $\mathcal{I}$ in $L^{2}(\Omega)$ is finite and less than or equal to $d$, provided $d$ is an integer number larger than $\max \left\{d_{1}, d_{2}\right\}$, where

$$
d_{1}:=\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right)
$$

and

$$
d_{2}:=\frac{2}{(1-\delta) \lambda_{1}}\left(\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right)\left(\frac{1-\delta}{2} \lambda_{1}-\mu_{1}\left(A_{\delta}\right)\right)+D\left(\gamma, \lambda_{0}, \delta,|\mathcal{I}|_{H^{1}}\right)\right) .
$$

Remark 4.9. The first proper value $\mu_{1}\left(A_{\delta}\right)$ of $A_{\delta}$ can be estimated from below in terms of $\lambda_{0}$ and $\left|\partial_{u} f(\cdot, 0)\right|_{L_{u}^{\sigma}}$. The explicit computations are left to the reader.

Remark 4.10. By Lemma 2.8, also the Hausdorff dimension of $\mathcal{I}$ in $H_{0}^{1}(\Omega)$ is finite and it is equal to the Hausdorff dimension of $\mathcal{I}$ in $L^{2}(\Omega)$.

## 5. Estimate of $\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right)$

In this section we shall obtain an explicit estimate for the number $\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right)$ in terms of the dominating potential $V_{\varepsilon}$ of Hypothesis 4.4. Our main tool is the celebrated Cwickel-Lieb-Rozenblum inequality, in its abstract formulation due to Rozenblum and Solomyak (see [21]). In order to exploit the CLR inequality, we need to make some assumption on the regularity of the open domain $\Omega$. Namely, we make the following assumption:

Hypothesis 5.1. The open set $\Omega$ is a uniformly $C^{2}$ domain in the sense of Browder [7, p. 36].

As a consequence, by elliptic regularity we have that

$$
D(-\Delta)=H^{2}(\Omega) \cap H_{0}^{1}(\Omega) \subset L^{\infty}(\Omega)
$$

In this situation, if $\omega \in L_{\mathrm{u}}^{\sigma}\left(\mathbb{R}^{3}\right)$ then the assignment $u \mapsto \omega u$ defines a relatively bounded perturbation of $-\Delta$ and therefore $D(-\Delta+\omega)=H^{2}(\Omega) \cap H_{0}^{1}(\Omega)$. It follows that $X^{\alpha} \subset L^{\infty}(\Omega)$ for $\alpha>3 / 4$ (see [11, Theorem 1.6.1]).

Set $\bar{\varepsilon}:=(1-\delta) \lambda_{1} / 4$. Define the bilinear forms

$$
\widetilde{a}_{\delta, \bar{\varepsilon}}(u, v):=(1-\delta)\left(\int_{\Omega} \nabla u \cdot \nabla v d x+\int_{\Omega} \beta u v d x\right)-3 \bar{\varepsilon} \int_{\Omega} u v d x
$$

for $u, v \in H_{0}^{1}(\Omega)$, and

$$
b_{\delta, \bar{\varepsilon}}(u, v):=-\int_{\Omega} V_{\bar{\varepsilon}} u v d x
$$

Moreover, set

$$
a_{\delta, \bar{\varepsilon}}(u, v):=\widetilde{a}_{\delta, \bar{\varepsilon}}(u, v)+b_{\delta, \bar{\varepsilon}}(u, v)
$$

and denote by $\widetilde{A}_{\delta, \bar{\varepsilon}}$ and $A_{\delta, \bar{\varepsilon}}$ the selfadjoint operators induced by $\widetilde{a}_{\delta, \bar{\varepsilon}}$ and $a_{\delta, \bar{\varepsilon}}$, respectively.

A simple computation shows that

$$
\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right) \leq n_{\delta, \bar{\varepsilon}}
$$

where $n_{\delta, \bar{\varepsilon}}$ is the number of negative eigenvalues of $A_{\delta, \bar{\varepsilon}}$.
By Theorem 1.3.2 in [8], the operator $\widetilde{A}_{\delta, \bar{\varepsilon}}$ is positive (with $\widetilde{A}_{\delta, \bar{\varepsilon}} \geq \bar{\varepsilon} I$ ) and order preserving. Moreover, since $D\left(A_{\delta, \bar{\varepsilon}}^{\alpha}\right) \subset L^{\infty}(\Omega)$ for $\alpha>3 / 4$, then for every such $\alpha$ and $\gamma<\bar{\varepsilon}$ we have

$$
\left|e^{-t \widetilde{A}_{\delta, \bar{\varepsilon}}} u\right|_{L^{\infty}} \leq M_{\alpha, \gamma} t^{-\alpha} e^{-\gamma t}|u|_{L^{2}}, \quad u \in L^{2}(\Omega)
$$

where $M_{\alpha, \gamma}$ is a constant depending only on $\alpha, \gamma$ and on the embedding constant of $H^{2}(\Omega)$ into $L^{\infty}(\Omega)$. It follows that

$$
M_{\widetilde{A}_{\delta, \bar{\varepsilon}}}(t):=\left\|e^{-(t / 2) \widetilde{A}_{\delta, \bar{\varepsilon}}}\right\|_{\mathcal{L}\left(L^{2}, L^{\infty}\right)}^{2} \leq M_{\alpha, \gamma}^{2} 2^{2 \alpha} t^{-2 \alpha} e^{-\gamma t}
$$

We are now in a position to apply Theorem 2.1 in [21]. We have thus proved the following theorem:

Theorem 5.2. Assume that Hypotheses 2.2, 2.6, 4.4 and 5.1 are satisfied. Let $\bar{\varepsilon}:=(1-\delta) \lambda_{1} / 4$. Then

$$
\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right) \leq n_{\delta, \bar{\varepsilon}} \leq C_{q / 2} M_{q / 2, \gamma} \int_{\Omega} V_{\bar{\varepsilon}}(x)^{q} d x
$$

where $C_{\alpha}$ is a constant depending only on $\alpha$, for $\alpha>3 / 4$.

## 6. Dissipative equations: dimension of the attractor

In this section we specialize our results to the case of a dissipative equation. We make the following assumption:

Hypothesis 6.1. There exists a non negative function $D \in L^{q}(\Omega), 2 \geq q>$ $3 / 2$, such that

$$
\begin{equation*}
f(x, u) u \leq D(x)|u|, \quad(x, u) \in \Omega \times \mathbb{R} \tag{6.1}
\end{equation*}
$$

Remark 6.2. Hypotheses 6.1 and 2.2 together are equivalent to the structure assumption of Theorem 4.4 in [4].

An easy computation shows that $|f(x, 0)| \leq D(x)$ for $x \in \Omega$, and that $F(x, u):=\int_{0}^{u} f(x, s) d s$ satisfies

$$
F(x, u) \leq D(x)|u|, \quad(x, u) \in \Omega \times \mathbb{R} .
$$

By slightly modifying some technical arguments in [17], one can prove that the semiflow $\pi$ generated by equation (2.2) in $H_{0}^{1}(\Omega)$ possesses a compact global attractor $\mathcal{A}$. Moreover, $\pi$ is gradient-like with respect to the Lyapunov functional

$$
\mathcal{L}(u):=\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} \beta(x)|u|^{2} d x-\int_{\Omega} F(x, u) d x, \quad u \in H^{1}(\Omega) .
$$

Assuming Hypothesis 6.1, we shall give an explicit estimate for $|\mathcal{A}|_{H^{1}}$ in terms of $|D|_{L^{q}}$. Moreover, we shall prove that Hypothesis 6.1 implies Hypothesis 4.4, and we explicitly compute the dominating potential $V_{\varepsilon}$ in terms of $D$. Therefore, we are able to obtain an explicit estimate for the number $\mathcal{N}\left(\delta, \frac{1-\delta}{2} \lambda_{1}\right)$ in terms of $|D|_{L^{q}}$. As a consequence, the estimate of the dimension of $\mathcal{A}$ given by Theorem 4.8 can be made completely explicit in terms of the structure parameters of equation (1.1).

We have the following theorem:

Theorem 6.3. Assume Hypotheses 2.2, 2.6 and 6.1 are satisfied.
(a) Let $\phi \in H_{0}^{1}(\Omega)$ be an equilibrium of $\pi$. Then

$$
|\phi|_{H^{1}} \leq \frac{M_{q^{\prime}}}{\lambda_{0}}|D|_{L^{q}}
$$

where $M_{q^{\prime}}$ is the embedding constant of $H_{0}^{1}\left(\mathbb{R}^{3}\right)$ into $L^{q^{\prime}}\left(\mathbb{R}^{3}\right)$.
(b) There exists a constant $S>0$ such that

$$
|u|_{H^{1}} \leq S \quad \text { for all } u \in \mathcal{A}
$$

The constant $S$ can be explicitly computed and depends only on $C, \gamma, \sigma$, $\lambda_{0}, \Lambda_{0},|D|_{L^{q}},\left|\partial_{u} f(\cdot, 0)\right|_{L_{\mathrm{u}}^{\sigma}}$ and on the constants of Sobolev embeddings.

Proof. Let $\phi \in H_{0}^{1}(\Omega)$ be an equilibrium of $\pi$. Then, for $\varepsilon>0$, we have

$$
\begin{aligned}
\lambda_{0}|\phi|_{H^{1}}^{2} & \leq \int_{\Omega}|\nabla \phi|^{2} d x+\int_{\Omega} \beta(x)|\phi|^{2} d x=\int_{\Omega} f(x, \phi) \phi d x \leq \int_{\Omega} D(x)|\phi| d x \\
& \leq|D|_{L^{q}}|\phi|_{L^{q^{\prime}}} \leq \varepsilon|\phi|_{L^{q^{\prime}}}^{2}+\frac{1}{4 \varepsilon}|D|_{L^{q}}^{2} \leq \varepsilon M_{q^{\prime}}^{2}|\phi|_{H^{1}}^{2}+\frac{1}{4 \varepsilon}|D|_{L^{q}}^{2}
\end{aligned}
$$

choosing $\varepsilon:=\lambda_{0} /\left(2 M_{q^{\prime}}^{2}\right)$ we get property (a). In order to prove (b), we notice that, since $\mathcal{L}$ is a Lyapunov functional for $\pi$ and $\mathcal{A}$ is compact in $H_{0}^{1}(\Omega)$, there exists an equilibrium $\phi$ such that, for every $u \in \mathcal{A}$,

$$
\begin{aligned}
\int_{\Omega}|\nabla u|^{2} d x+\int_{\Omega} \beta(x)|u|^{2} d x & -\int_{\Omega} F(x, u) d x \\
& \leq \int_{\Omega}|\nabla \phi|^{2} d x+\int_{\Omega} \beta(x)|\phi|^{2} d x-\int_{\Omega} F(x, \phi) d x
\end{aligned}
$$

Then, for $\varepsilon>0$, we have:

$$
\begin{aligned}
\lambda_{0}|u|_{H^{1}}^{2} & \leq \int_{\Omega} D(x)|u| d x+\Lambda_{0}|\phi|_{H^{1}}^{2}+\int_{\Omega} F(x, \phi) d x \\
& \leq \varepsilon M_{q^{\prime}}^{2}|u|_{H^{1}}^{2}+\frac{1}{4 \varepsilon}|D|_{L^{q}}^{2}+\Lambda_{0}|\phi|_{H^{1}}^{2}+\int_{\Omega} F(x, \phi) d x
\end{aligned}
$$

We choose $\varepsilon:=\lambda_{0} /\left(2 M_{q^{\prime}}^{2}\right)$ and the conclusion follows.
Finally, we have:
Theorem 6.4. Assume that Hypotheses 2.6 and 6.1 are satisfied. Then for every $0<\varepsilon \leq 1$,

$$
\partial_{u} f(x, 0) \leq \frac{2}{\varepsilon} D(x)+\frac{\varepsilon}{2} C\left(1+\varepsilon^{\gamma}\right)
$$

Proof. For $\varepsilon>0$ we have:

$$
f(x, \varepsilon)=f(x, 0)+\partial_{u} f(x, 0) \varepsilon+\int_{0}^{\varepsilon}\left(\int_{0}^{s} \partial_{u u} f(x, r) d r\right) d s
$$

It follows that

$$
f(x, 0) \varepsilon+\partial_{u} f(x, 0) \varepsilon^{2}+\varepsilon \int_{0}^{\varepsilon}\left(\int_{0}^{s} \partial_{u u} f(x, r) d r\right) d s=f(x, \varepsilon) \varepsilon \leq D(x) \varepsilon
$$

Therefore

$$
\partial_{u} f(x, 0) \leq \frac{D(x)+|f(x, 0)|}{\varepsilon}+\frac{1}{\varepsilon} \int_{0}^{\varepsilon}\left(\int_{0}^{s} C\left(1+|r|^{\gamma}\right) d r\right) d s
$$

and the conclusion follows.
Remark 6.5. Theorem 6.4 shows that Hypotheses 2.6 and 6.1 together imply Hypothesis 4.4 , with $V_{\varepsilon}(x)=\frac{2 C}{\varepsilon} D(x)$.

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[^0]:    2010 Mathematics Subject Classification. 35L70, 35B40, 35B65.
    Key words and phrases. Reaction diffusion equation, invariant set, attractor, dimension

